

Measurement of Semileptonic B Decays into Orbitally Excited Charmed Mesons

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We present a study of B decays into semileptonic final states containing charged and neutral $D_1(2420)$ and $D_2^*(2460)$. The analysis is based on a data sample of 208 fb^{-1} collected at the $Y(4S)$ resonance with the $BABAR$ detector at the PEP-II asymmetric-energy B factory at SLAC. With a simultaneous fit to four different decay chains, the semileptonic branching fractions are extracted from measurements of the mass difference $\Delta m = m(D^{**}) - m(D)$ distributions. Product branching fractions are determined to be $\mathcal{B}(B^+ \rightarrow D_1^0 \ell^+ \nu_\ell) \times \mathcal{B}(D_1^0 \rightarrow D^{*+} \pi^-) = (2.97 \pm 0.17 \pm 0.17) \times 10^{-3}$, $\mathcal{B}(B^+ \rightarrow D_2^{*0} \ell^+ \nu_\ell) \times \mathcal{B}(D_2^{*0} \rightarrow D^{(*)+} \pi^-) = (2.29 \pm 0.23 \pm 0.21) \times 10^{-3}$, $\mathcal{B}(B^0 \rightarrow D_1^- \ell^+ \nu_\ell) \times \mathcal{B}(D_1^- \rightarrow D^{*0} \pi^-) = (2.78 \pm 0.24 \pm 0.25) \times 10^{-3}$ and $\mathcal{B}(B^0 \rightarrow D_2^{*-} \ell^+ \nu_\ell) \times \mathcal{B}(D_2^{*-} \rightarrow D^{(*)0} \pi^-) = (1.77 \pm 0.26 \pm 0.11) \times 10^{-3}$. In addition we measure the branching ratio $\Gamma(D_2^* \rightarrow D \pi^-) / \Gamma(D_2^* \rightarrow D^{(*)} \pi^-) = 0.62 \pm 0.03 \pm 0.02$.

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Measurements of the Cabbibo-Kobayashi-Maskawa matrix elements $|V_{cb}|$ and $|V_{ub}|$ rely on precise knowledge of semileptonic B -meson decays. Decays with orbitally-excited charm mesons (D^{**}) in the final state give a significant contribution to the total semileptonic decay rate. A better understanding of these decays will reduce the uncertainty in the composition of the signal and backgrounds for inclusive and exclusive measurements [1].

In the framework of heavy quark symmetry (HQS), D^{**} mesons form two doublets with $j_q^P = 1/2^-$ and $j_q^P = 3/2^-$ where j_q^P denotes the spin-parity of the light quark coupled to the orbital angular momentum. The doublets with $j_q^P = 3/2^-$, namely, the D_1 and D_2^* , have to decay via D wave to conserve parity and angular momentum and therefore are narrow with widths of order of 10 MeV [2]. The relative contribution of the two doublets and the polarization of the produced D^{**} mesons can be compared with QCD sum rules [3] and predictions from heavy quark effective theory [4].

In this Letter we describe a simultaneous measurement of all B semileptonic decays to the two narrow orbitally-excited charmed states, without explicit reconstruction of the rest of the event. The CLEO collaboration has previously reported a branching fraction measurement for $B^+ \rightarrow D_1^0 \ell^+ \nu$ and an upper limit for $B^+ \rightarrow D_2^{*0} \ell^+ \nu$ [5]. Belle and $BABAR$ have reported results using a technique in which one of the B mesons in the process $Y(4S) \rightarrow B\bar{B}$ is fully reconstructed [6].

In this analysis we use a sample with a total integrated luminosity of 208 fb^{-1} , part of the complete data set collected with the $BABAR$ detector at the PEP-II storage ring, operating at a center of mass energy of 10.58 GeV.

The $BABAR$ detector [7] and event reconstruction [8] are described in detail elsewhere. A Monte Carlo (MC) simulation of the detector based on GEANT4 [9] is used to estimate signal efficiencies and to understand the backgrounds. The sample of simulated $B\bar{B}$ events is equivalent to approximately 3 times the data sample and a dedicated simulation of signal events based on the ISGW2 model [10] has been produced with statistics equivalent to

roughly 5 times the expected signal yield contained in the data.

D^{**} decays are reconstructed in the decay chains $D^{**} \rightarrow D^* \pi^-$ [11], and $D^{**} \rightarrow D \pi^-$. The former is accessible to both narrow D^{**} states while the latter has no contribution from the D_1 . Intermediate D^* states are reconstructed in $D^* \rightarrow D^0 \pi$ and the D mesons are reconstructed exclusively in $D^0 \rightarrow K^- \pi^+$ and $D^+ \rightarrow K^- \pi^+ \pi^+$. D^{**} candidates are then paired with reconstructed leptons and required to be consistent with the semileptonic decays $B \rightarrow D^{**} \ell \nu$, as described in the following.

First, events which are most likely to contain a semileptonic B decay are selected. We require that there is a reconstructed D candidate and at least one lepton in the event with a momentum greater than 800 MeV/c [12]. D^0 meson candidates are formed by $K^- \pi^+$ combinations requiring the invariant mass to be consistent with the D^0 mass: $1.846 < m(K\pi) < 1.877 \text{ GeV}/c^2$. This asymmetric mass window is chosen to take into account resolution effects of the detector. The selection is optimized to maximize the significance of the selected sample.

D^0 candidates are combined with charged and neutral pions to form D^* candidates. For D^{*0} the π^0 is reconstructed from a photon pair with an invariant mass of $115 < m_{\gamma\gamma} < 150 \text{ MeV}/c^2$. Those photon pairs are refitted in a “mass-constrained” fit to match the nominal mass of the π^0 . D^* candidates are selected by their mass difference to the D^0 candidate: $144 < m(D^0 \pi^+) - m(D^0) < 148 \text{ MeV}/c^2$ and $140 < m(D^0 \pi^0) - m(D^0) < 144 \text{ MeV}/c^2$ for charged and neutral D^* , respectively.

D^+ candidates are formed from $K^- \pi^+ \pi^+$ combinations with an invariant mass of $1.854 < m(K\pi\pi) < 1.884 \text{ GeV}/c^2$. The χ^2 fit probability for the three tracks to originate from a common vertex, P_{vtx} , is required to be $P_{\text{vtx}}(K\pi\pi) > 0.01$.

Candidates for D and D^* are combined with charged pions to form D^{**} candidates, and finally paired with muons or electrons. The charge of the lepton is required to match the charge of the kaon from the D decay.

Part of the background is due to events where a D^{**} is paired to a lepton from the other B . Thus we require that the probability that the lepton and the pion emitted by the D^{**} originate from a common vertex exceeds 0.001, and that the angle between the direction of flight of the D^{**} and the lepton is more than 90 degrees.

A large fraction of the background events is due to $B \rightarrow D^* \ell \nu$ decays where the D^* or its daughter D is paired to a pion from the other B . To suppress this combinatorial background, we make use of the variable \cos_{BY} described in the following. The energy and momentum of the B mesons from the $Y(4S)$ decays are known from incident beam energies. For correctly reconstructed $B \rightarrow D^{**} \ell \nu$ decays, where the only missing particle is the neutrino, the decay kinematics can be calculated, up to one angular quantity, from the four-momentum of the visible decay products ($Y = D^{**} \ell$). The cosine of the angle between the direction of flight of the B meson and its visible decay product Y is given by

$$\cos_{BY} = -\frac{2E_B E_Y - m_B^2 - m_Y^2}{2|\vec{p}_B||\vec{p}_Y|},$$

where E , $|\vec{p}|$ and m are the energies, momenta, and masses of the B and the Y , respectively. If the Y candidate is not from a correctly reconstructed $B \rightarrow D^{**} \ell \nu$ decay, the quantity \cos_{BY} no longer represents an angle, and can take any value. We select candidates having $|\cos_{BY}| \leq 1$.

In case a D^* is reconstructed in the decay chain, a veto is applied against decays $B \rightarrow D^* \ell \nu$ by calculating the variable $\cos_{BY'}$ which is defined as above, but the Y system is redefined to contain only the D^* and the lepton: $Y' = D^* \ell$. Background events are rejected by the requirement $\cos_{BY'} < -1$ since signal events $B \rightarrow D^{**} \ell \nu$ tend to have values less than -1 .

To reduce combinatorial backgrounds in the decay chain $D^{**} \rightarrow D \pi^-$, only the $D^{**} \ell$ candidate with \tilde{m}_ν^2 closest to zero is selected, where \tilde{m}_ν^2 is the neutrino mass squared, calculated in the approximation $\vec{p}_B = 0$: $\tilde{m}_\nu^2 = m_B^2 + |\vec{p}_Y|^2 - 2E_B E_Y$. Events reconstructed in the $D^{**} \rightarrow D^0 \pi^-$ final state are rejected if the D^0 can be paired with any charged pion to form a D^{*+} candidate as described above.

In about 2% of the events more than one $D^{**} \ell$ candidate is selected and if so all of them enter the analysis.

We determine the D_2^* signal yield in the channel $D^{**} \rightarrow D \pi$ and the D_1 and D_2^* signal yields in the channel $D^{**} \rightarrow D^* \pi$ by a binned χ^2 fit to the $\Delta m = m(D^{(*)} \pi^-) - m(D^0)$ distributions. To determine the individual contributions from D_1 and D_2^* in the $D^* \pi$ final state, we make use of the helicity angle distribution of the D^* , ϑ_h , which is defined as the angle between the two pions emitted by the D^{**} and the D^* in the rest frame of the D^* . For a D^* from a D_2^* this distribution varies as $\sin^2 \vartheta_h$, whereas for D_1 decays, the helicity angle is distributed like $1 +$

$A_{D_1} \cos^2 \vartheta_h$, where A_{D_1} is a parameter which depends on the initial polarization of the D_1 and a possible S -wave contribution to the D_1 decay. To exploit this feature, we split the data for the two decay chains involving a D^* into four subsamples, corresponding to four equal size bins in $|\cos \vartheta_h|$.

The resulting ten Δm distributions are fitted simultaneously to determine 12 parameters describing the signal yields and distributions, and 22 parameters to adjust the background yields and shapes. The mass differences for the signal events are described by Breit-Wigner functions. There are four parameters giving the signal yields for the semileptonic decays involving the two narrow states, charged and neutral. The masses of the states are also fitted, but are constrained to be equal for charged and neutral states, giving two parameters. Four additional parameters arise from the effective widths of the D^{**} states, which represent a convolution of the intrinsic widths and detector resolution effects. The latter contributes approximately $2-3 \text{ MeV}/c^2$, depending on the mode. The fit also determines the D_2^* branching ratio $\mathcal{B}_{D/D^{(*)}} = \Gamma(D_2^* \rightarrow D \pi^-) / (\Gamma(D_2^* \rightarrow D \pi^-) + \Gamma(D_2^* \rightarrow D^* \pi^-))$ and the D_1 polarization amplitude A_{D_1} .

Backgrounds are modeled by cubic functions in Δm . The background shape in the $D^* \pi^-$ channel is found to be the same in all helicity bins for each final state. The fit thus has three shape parameters for each decay chain, while the number of background events is determined independently in each bin.

The selection efficiency is deduced from a fit to the simulation. This fit uses the same parametrization as the fit determining the signal yield from data and is applied to the sum of the full background simulation and for one signal decay chain at a time. For a given decay mode the efficiencies are found to be the same for D_1 and D_2^* , specifically: $\epsilon(D^{*+} \pi^-) = (6.89 \pm 0.12)\%$, $\epsilon(D^{*0} \pi^-) = (5.34 \pm 0.12)\%$, $\epsilon(D^+ \pi^-) = (12.88 \pm 0.96)\%$ and $\epsilon(D^0 \pi^-) = (17.56 \pm 0.70)\%$, where the quoted uncertainties are the statistical uncertainties from the fit. For the decays including a D^* the efficiency is multiplied by the probability for a D^{**} to decay with a value of $|\cos \vartheta_h|$ falling into a given bin. This factor includes the theoretical distribution discussed above as well as corrections for the different detector acceptances in the four helicity bins of up to 10%. The total number of B mesons in the data sample used for the present work is $N_{B\bar{B}} = (236.0 \pm 2.6) \times 10^6$ [13]. For the charged and neutral B mesons we assume $\Gamma(Y(4S) \rightarrow B^+ B^-) / \Gamma(Y(4S) \rightarrow B^0 \bar{B}^0) = 1.065 \pm 0.026$ [14].

The fit procedure has been extensively validated. The analysis procedure is tested on statistically independent MC simulated data samples and was found to reproduce the input signal parameters with a $\chi^2/n = 12.66/12$, where n is the number of signal parameters. Consistent fit results were also obtained when the data sample was

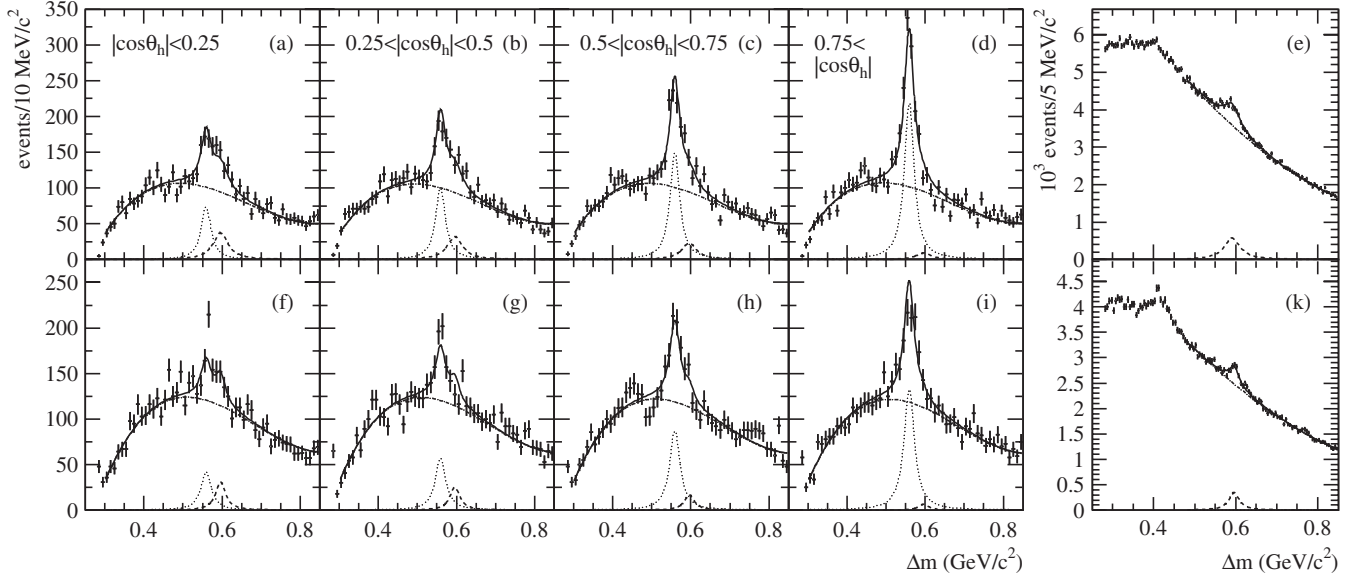


FIG. 1. Δm spectra for the selected data and the results of the fitted functions. The solid line represents the complete fit function, dotted (D_1) and dashed (D_2^*) lines for the signal and dash-dotted the for background. (a) to (d) show the mode $D^{**0} \rightarrow D^{**+} \pi^-$ with increasing values for $|\cos \vartheta_h|$, (e) the mode $D^{**0} \rightarrow D^+ \pi^-$. (f) to (i) show the corresponding bins in $|\cos \vartheta_h|$ for the mode $D^{**+} \rightarrow D^{*0} \pi^+$ and (k) the mode $D^{**+} \rightarrow D^0 \pi^+$.

separated into subsamples representing specific data taking periods, separated by lepton species or restricting it to certain decay modes, using charged or neutral D^{**} only, or combining the helicity bins. The results of the fit are shown in Fig. 1. As expected, the contribution of the D_2^* vanishes for large values of $|\cos \vartheta_h|$ while the contribution of the D_1 is suppressed for $\cos \vartheta_h$ close to zero. The extracted yields are given in Table I.

Systematic uncertainties have been analyzed and their impact on the fitted yields have been estimated taking into account correlations between fit parameters. Efficiencies for reconstructing and selecting the particles of the final state are derived from Monte Carlo simulation. The simulation of the tracking and the π^0 reconstruction have been studied by comparing τ decays to one and three charged tracks and with or without a neutral pion. Uncertainties introduced by the particle identification for kaons and leptons are studied using control samples with high purities for the particles in question. The impact of the finite statistics of the simulated signal events is deduced from the fit error of the efficiency determination.

TABLE I. Extracted yields for the four signal modes in the five relevant Δm spectra.

Mode	$ \cos \vartheta_h $	D_1^0	D_2^{*0}	D_1^+	D_2^{*+}
$D^* \pi^+$	[0.00 0.25]	344	273	212	152
$D^* \pi^+$	[0.25 0.50]	470	238	286	123
$D^* \pi^+$	[0.50 0.75]	699	170	439	83
$D^* \pi^+$	[0.75 1.00]	1027	67	668	31
$D \pi^+$	8414	...	3361

The uncertainty on the number of charged and neutral B mesons in the data set is determined as in [13,14] and the branching fractions of the decays of the D^* and the D are taken from [15].

Uncertainties introduced by the physics model which was used to simulate the MC data have been addressed by reweighting the signal MC calculations to an alternative decay model based on HQET [4]. The fit was repeated with efficiencies deduced from the reweighted signal MC data and the deviations in the results are taken as systematic uncertainties. A possible influence of the background description has been tested by varying the parametrizations.

TABLE II. Summary of systematic uncertainties of the determination of the semileptonic branching fractions.

Source	$\Delta \mathcal{B}(B \rightarrow D^{**} \ell \nu) / \mathcal{B}(B \rightarrow D^{**} \ell \nu) [\%]$			
	D_1^0	D_2^{*0}	D_1^+	D_2^{*+}
Tracking	1.76	1.39	1.03	1.14
π^0 efficiency	0.06	0.29	3.25	0.60
Particle identification	2.61	2.75	3.11	1.60
MC statistics	1.80	5.61	2.50	3.32
Helicity correction	0.65	0.14	0.17	0.31
Number of B mesons	2.68	2.68	2.68	2.68
$\mathcal{B}(D^{*+} \rightarrow D^0 \pi^+)$	0.76	0.19	0.04	0.10
$\mathcal{B}(D^{*0} \rightarrow D^0 \pi^0)$	0.11	0.45	5.07	0.93
$\mathcal{B}(D^0 \rightarrow K^- \pi^+)$	1.89	0.42	1.78	2.03
$\mathcal{B}(D^+ \rightarrow K^- \pi^+ \pi^+)$	0.07	2.67	0.24	0.54
Signal modeling	2.11	4.75	3.21	1.95
bkg. parametrization	1.93	1.68	3.20	2.71
Total	5.76	9.03	9.16	6.17

The backgrounds are alternatively described by a square root function, $f(\Delta m) = \sqrt{\Delta m - m_0}$, where m_0 is the kinematic limit, multiplied by either polynomials or exponentials in Δm .

Table II gives a summary of the various sources of systematic uncertainty and their impact on the results.

$$\begin{aligned} \mathcal{B}(B^+ \rightarrow D_1^0 \ell^+ \nu_\ell) \times \mathcal{B}(D_1^0 \rightarrow D^{*+} \pi^-) &= (2.97 \pm 0.17_{\text{stat}} \pm 0.17_{\text{syst}}) \times 10^{-3}, \\ \mathcal{B}(B^+ \rightarrow D_2^{*0} \ell^+ \nu_\ell) \times \mathcal{B}(D_2^{*0} \rightarrow D^{(*)+} \pi^-) &= (2.29 \pm 0.23_{\text{stat}} \pm 0.21_{\text{syst}}) \times 10^{-3}, \\ \mathcal{B}(B^0 \rightarrow D_1^- \ell^+ \nu_\ell) \times \mathcal{B}(D_1^- \rightarrow D^{*0} \pi^-) &= (2.78 \pm 0.24_{\text{stat}} \pm 0.25_{\text{syst}}) \times 10^{-3}, \\ \mathcal{B}(B^0 \rightarrow D_2^{*-} \ell^+ \nu_\ell) \times \mathcal{B}(D_2^{*-} \rightarrow D^{(*)0} \pi^-) &= (1.77 \pm 0.26_{\text{stat}} \pm 0.11_{\text{syst}}) \times 10^{-3}. \end{aligned}$$

We observe all modes with significance greater than 5σ , among them evidence of the D_2^{*-} contribution to the decay $B \rightarrow D^* \pi \ell \nu$. For modes already observed we find results in agreement with previous measurements, but achieve better precisions [5,6,16].

For the decays of the D^{**} we measure the branching ratio $\mathcal{B}_{D/D^{(*)}} = 0.62 \pm 0.03_{\text{stat}} \pm 0.02_{\text{syst}}$. This ratio is in agreement with theoretical predictions [2] and previous measurements [15] but has a smaller uncertainty by a factor of four.

For the D_1 we determine the polarization parameter to be $A_{D_1} = 3.8 \pm 0.6_{\text{stat}} \pm 0.8_{\text{syst}}$. It is the first measurement of the D_1 polarization, within the uncertainties consistent with unpolarized D_1 decaying purely via D wave, which gives the prediction $A_{D_1} = 3$, but violates HQS [4].

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Added in quadrature the total systematic uncertainties in the semileptonic branching fractions are 6%–10%, depending on the D^{**} type.

In summary, we have measured the four branching fractions of B mesons decaying semileptonically into narrow D^{**} states. The D^{**} decay rates are unknown; thus, we can only determine the product branching fractions:

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