

1 **Field-Induced SMM and Visible/NIR-Luminescence Behaviour of Dinuclear LnIII Complexes**  
2 **with 2-Fluorobenzoate**

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35 **ABSTRACT:**

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37 The reaction of  $\text{Ln}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$  salts ( $\text{Ln} = \text{Nd, Eu, Gd, Tb, Dy, Er}$  and  $\text{Yb}$ ) with 2-fluorobenzoic acid  
38 ( $\text{H-2-FBz}$ ) and 1,10-phenanthroline ( $\text{phen}$ ) in ethanol/water mixture allows the isolation of dinuclear  
39 compounds of the formula  $[\text{Ln}_2(\text{2-FBz})_4(\text{NO}_3)_2(\text{phen})_2]$   $\{\text{Ln} = \text{Nd}$  (1),  $\text{Eu}$  (2),  $\text{Gd}$  (3),  $\text{Tb}$  (4),  $\text{Dy}$  (5),  
40  $\text{Er}$  (6) $\}$  and  $[\text{Yb}_2(\text{2-FBz})_6(\text{phen})_2]$  (7). The solid-state photoluminescence study of the complexes  
41 shows the 4f–4f lanthanide transitions in the visible range, in the cases of 2, 4 and 5, and in the NIR  
42 range for 1, 6 and 7. Magnetic studies reveal field-induced single-molecule-magnet (SMM) behaviour  
43 for compounds 1, 5, 6 and 7.

44

## 45 INTRODUCTION

46

47 Multifunctional molecular materials can be defined as compounds that involve coexistence, interplay or  
48 synergy between multiple physical properties.[1,2] Restricting the discussion to the combination of  
49 luminescence and single-molecule magnet (SMM) properties, the LnIII coordination compounds are  
50 ideal candidates for the construction of this kind of hybrid molecular material.[3–8] On the one hand,  
51 the 4f–4f electronic transitions (responsible for the light emission) are narrow and characteristic for each  
52 LnIII, and the emitting excited states are long-lived. Since the f–f transitions are parity-forbidden, free  
53 LnIII ions have low extinction coefficients, leading to low luminescence intensity. To overcome this  
54 problem, the presence of light-harvesting ligands coordinated to the LnIII can enhance the metal  
55 luminescence through an energy transfer process, commonly known as the antenna effect.[9] The R–  
56 benzoate ligands have been widely used in Ln coordination compounds due to the strong interaction  
57 between the LnIII ions and the oxygen atoms from the carboxylate group and also because of the strong  
58 absorbing chromophore group of the organic fragment.[10] The use of a fluorinated benzoate, such as 2-  
59 fluorobenzoic acid, could enhance the NIR emission of Nd, Er and Yb, since the C–F vibrational  
60 quenching is much lower than the one produced by the C–H bonds.[11] Moreover, chelating ligands  
61 such as 1,10-phenanthroline (phen) can block two coordination sites per LnIII ion and can terminate  
62 further aggregation or potential polymerization.[12] The 1,10-phenanthroline ligand can also sensitize  
63 the luminescence of lanthanide ions through their large pi system.[13]

64 On the other hand, some compounds containing LnIII ions with high anisotropic magnetic moments  
65 show SMM properties.[14] Among the lanthanide ions, DyIII has yielded the largest number of 4f-based  
66 SMMs, followed by TbIII.[9,15,16,17] Nevertheless, ErIII and YbIII are also good candidates to present  
67 SMM properties.[18–20] Slow relaxation of the magnetization has been observed on several lanthanide  
68 compounds containing different R–benzoate ligands.[21–25]

69 In this work, we present the structure of six new dinuclear 4f-metal complexes by simultaneously using  
70 bidentate bridging carboxylate groups derived from 2-fluorobenzoic acid (H-2-FBz), chelating 1,10-  
71 phenanthroline (phen) and nitrate capping ligands. Four of the new reported compounds have the  
72 formula  $[\text{Ln}_2(2\text{-FBz})_4(\text{NO}_3)_2(\text{phen})_2]$  {Ln = Nd (1), Eu (2), Ln = Gd (3), Ln = Dy (5), Ln = Er (6)} and  
73 are isostructural to the previously reported[26]  $[\text{Tb}_2(2\text{-FBz})_4(\text{NO}_3)_2(\text{phen})_2]$  (4) compound. The  
74 coordination number of the lanthanide ions is 9 for 1–6. The same synthetic procedure for Ln = Yb  
75 yielded a new dinuclear compound, without the nitrate ligand, with formula  $[\text{Yb}_2(2\text{-FBz})_6(\text{phen})_2]$  (7).  
76 The coordination number of the YbIII ion in this compound is 8. For the new compounds, we report  
77 here the syntheses, crystal structure, magnetic behaviour and luminescence properties. For the  
78 previously published  $[\text{Tb}_2(2\text{-FBz})_4(\text{NO}_3)_2(\text{phen})_2]$  (4) compound, as only the structure was reported,  
79 we have also studied its magnetic and luminescence properties.

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## 82 RESULTS AND DISCUSSION

83

### 84 X-ray Diffraction Crystal Structures

85 Complexes 1–6 are isomorphs and crystallize in the triclinic space group  $P1\bar{1}$ ; thus, as an example, only  
86 the structure of 2 will be discussed in detail.

87

#### 88 [Eu<sub>2</sub>(2-FBz)<sub>4</sub>(NO<sub>3</sub>)<sub>2</sub>(phen)<sub>2</sub>] (2)

89 A partially labelled plot of the structure of the dinuclear compound 2 is shown in Figure 1a. Selected  
90 bond lengths are listed in Table 1. The structure of compound 2 consists of centrosymmetric dinuclear  
91 molecules in which each EuIII is nine-coordinate. The two EuIII atoms are bridged through four 2-  
92 fluorobenzoate ligands with two different kinds of bridging modes. One of them is a symmetrical syn-  
93 syn bridge ( $\eta^1:\eta^1:\mu$  or 2.11 using Harris notation) (Scheme 1a), with the Eu1–O1 and Eu1–O2' bond  
94 lengths being 2.3698(1) and 2.3531(1) Å, respectively. The second type of 2-fluorobenzoate bridging  
95 ligands are best described as chelating–bridging ( $\eta^1:\eta^2:\mu$  or 2.21) (Scheme 1b), in which O3 acts as a  
96 bridge between the two Eu atoms, with a distance of 2.6989(1) for Eu1–O3 and 2.3705(1) Å for Eu1–  
97 O3'; meanwhile, O4 is bonded only to one Eu with a 2.4066(1) Å bond length. The intramolecular  
98 Eu1···Eu1' distance is 3.9605(1) Å. The coordination sphere of the metals is completed by the N1 and  
99 N2 atoms of a phenanthroline ligand with Eu1–N1 and Eu1–N2 bond lengths of 2.5811(1) and  
100 2.6170(1) Å, respectively, and by the O5 and O6 atoms from a nitrate anion with Eu1–O5 and Eu1–O6  
101 bond lengths of 2.5329(1) and 2.4604(1) Å, respectively. The calculation of the degree of distortion of  
102 the EuN<sub>2</sub>O<sub>7</sub> coordination polyhedron for 2 with respect to the ideal nine-vertex polyhedron, by using  
103 continuous shape measurement theory and SHAPE software[27] shows that the EuN<sub>2</sub>O<sub>7</sub> arrangement  
104 for 2 is intermediate between various coordination polyhedra (Table S1 in the Supporting Information).  
105 The best SHAPE calculations lead to Muffin (MFF-9), spherical capped square antiprism (CSAPR-9)  
106 and tricapped trigonal prism (JTCTPR-9) geometries, with continuous shape measure (CShM) values of  
107 2.170, 2.327 and 2.819, respectively. The calculations of the degree of distortion of the LnN<sub>2</sub>O<sub>7</sub>  
108 coordination polyhedra for 1, 3, 5 and 6, with respect to the ideal nine-vertex polyhedron, are also  
109 shown in Table S1. A graphical representation of the metal coordination geometry of 2 is shown in  
110 Figure 1b.

111 The dinuclear entities are assembled into 1D chains in the [101] direction through  $\pi$ -stacking  
112 interactions between the central rings of two adjacent phenanthroline ligands from two different units  
113 (Figure 2). These chains are connected, forming a 2D sheet in the (111) plane through another  $\pi$ -  
114 stacking bond between two adjacent phenyl rings from the 2-F-benzoate chelating/bridging ligand. No  
115 classical hydrogen bonds are found in the crystal structures of complexes 1–3, 5 and 6, but there is an  
116 intramolecular hydrogen bond formed by C26–H26···F1, with a distance of 3.115(3) Å and an angle of  
117 136° (Figure 2). The  $\pi$ -stacking intermolecular interaction distances and intramolecular hydrogen bond  
118 lengths and angles for compounds 1–3, 5 and 6 are summarized in Tables S2 and S3.

## 119 **[Yb<sub>2</sub>(2-FBz)<sub>6</sub>(phen)<sub>2</sub>] (7)**

120 Complex 7 also crystallizes in the triclinic space group  $P1\bar{1}$  and consists of a centrosymmetric dinuclear  
121 unit, but in this case, each YbIII centre presents an octacoordinate environment. A partially labelled  
122 structure of the dinuclear compound 7 is shown in Figure 3a. Selected bond lengths are listed in Table 1.  
123 The two equivalent YbIII ions are connected through four oxygen atoms from two 2-F-benzoate  
124 bridging ligands in a syn–syn coordination mode (Scheme 1a), with an intramolecular Yb1⋯Yb1'  
125 distance of 5.152 Å. The Yb–O1 and Yb1–O2' distances are 2.262(2) and 2.206(2) Å, respectively. Each  
126 YbIII ion is bonded to two chelating 2-F-benzoates (Scheme 1c), with the Yb1–O3, Yb1–O4, Yb1–O5  
127 and Yb1–O6 distances ranging from 2.339 to 2.409 Å. The two nitrogen atoms of a phenanthroline  
128 ligand complete the coordination sphere of each ion, featuring Yb1–N1 and Yb1–N2 bond lengths of  
129 2.459(3) and 2.385(6) Å, respectively. The calculation of the degree of distortion of the YbN<sub>2</sub>O<sub>6</sub>  
130 coordination polyhedron for 7, with respect to the ideal eight-vertex polyhedron, by using the  
131 continuous shape measure theory and SHAPE software,[20] shows that the YbN<sub>2</sub>O<sub>6</sub> arrangement for 7  
132 is intermediate between various coordination polyhedra. The best SHAPE calculations lead to  
133 biaugmented trigonal prism (BTPR-8), triangular dodecahedron (TDD-8) and Johnson biaugmented  
134 trigonal prism (JBTPR-8) geometries, with CShM values of 2.551, 3.135 and 3.145, respectively. A  
135 graphical representation of the metal coordination geometry for 7 is shown in Figure 3b.  
136 In complex 7, there are no classical hydrogen bonds. Nevertheless, there is a weak intermolecular  
137 interaction between C32–H32 with a F1 atom of an adjacent dinuclear fragment that leads to the  
138 arrangement of these molecules into a 1D chain along the [011] direction, which is enhanced by a  $\pi$ -  
139 stacking interaction between two phenanthroline ligands, as is depicted in Figure 4. Moreover, the  
140 chains are connected, giving a 2D sheet in the (111) plane through the contact between C12–H12 and F3  
141 from another unit (Figure 4). The intermolecular  $\pi$ -stacking interaction distance and intra- and  
142 intermolecular hydrogen-bond lengths and angles are summarized in Tables S2 and S3.

143

## 144 **Synthesis**

145 To avoid the highly energetic and not very controllable hydroor solvothermal processes, in this work, we  
146 used a straightforward room-temperature synthetic procedure, different from that used to prepare the  
147 already published complex [Tb<sub>2</sub>(2-FBz)<sub>4</sub>(NO<sub>3</sub>)<sub>2</sub>(phen)<sub>2</sub>],[26] and have successfully obtained the latter  
148 and six new dinuclear lanthanide compounds. Five of the mentioned complexes are isostructural with  
149 the TbIII one, but when using YbIII as the lanthanide source, the complex obtained is different and has  
150 the formula [Yb<sub>2</sub>(2-FBz)<sub>6</sub>(phen)<sub>2</sub>].

151 It is worth noting, then, the viability of the room-temperature approach for obtaining coordination  
152 lanthanide compounds.

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## 156 **Structural Discussion**

157 This family of  $[\text{Ln}_2(2\text{-FBz})_4(\text{NO}_3)_2(\text{phen})_2]$  complexes provides an opportunity to study the influence  
158 of the lanthanide contraction over the structural arrangement. When the atomic number of the metal  
159 increases, the radius of the LnIII cation decreases, and slight changes in the coordination sphere of the  
160 Ln can be observed, as has been demonstrated in other papers.[28–30] The most important change  
161 observed herein is the two different types of structures obtained in this work by using the same synthetic  
162 approach. For NdIII, EuIII, GdIII, TbIII, DyIII and ErIII ions, a homo-dinuclear structure was obtained,  
163 in which the nonacoordinated metals are bridged by four carboxylate ligands, two in syn–syn  
164 coordination mode and two in chelating/bridging mode. Meanwhile, in the case of the dinuclear  
165 complex with the YbIII ion, with the smallest radius of the family, the coordination decreases to eight,  
166 and therefore, the two metals inside the molecule are bridged by only two carboxylate ligands in a syn–  
167 syn coordination mode. Additionally, there is a terminal carboxylate, instead of the nitrate anion, present  
168 in the former structural type.

169 Other systematic structural variations coming from the different size of the ionic radius can be extracted  
170 for each type of structure (Table 1). Within the isostructural complexes 1–3, 5, 6 and the previously  
171 published compound 4, the bond lengths from the coordination sphere of the LnIII decrease 2–3 % from  
172 Eu to Er, except for the largest Ln1–O3 bond, which presents a tendency to increase nearly 2 % with the  
173 diminution of the LnIII radius. Then, for almost all of the Ln–O and Ln–N bonds, the lengths decrease  
174 due to the increase of the lanthanide contraction along the period.

175 On the other hand, the intramolecular LnIII···LnIII distances are almost constant in the entire studied  
176 series for 1–6, with an average value of 3.960 Å. Due to the different coordination mode of the 2-FBz  
177 bridging ligands in compound 7, the YbIII···YbIII distance is 5.152 Å, and it provides the largest Ln–Ln  
178 distance between the metal atoms of all of the compounds presented in this work.

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180

## 181 **Magnetic Properties**

182

### 183 **Dc Magnetic Susceptibility Study**

184 Solid-state direct-current (dc) magnetic susceptibility ( $\chi_M$ ) data on polycrystalline powder samples of  
185 complexes 1–7 were collected under applied magnetic fields of 0.3 T (300–2 K) for 1 and 3–7 and at 0.5  
186 T (300–2 K) for 2. The data are plotted as  $\chi_M \times T$  versus T in Figure 5.

187 All of the compounds presented in this work are homodinuclear carboxylate-bridged lanthanide  
188 compounds. By analogy with the dn–dn dinuclear compounds, the exchange coupling interaction can be  
189 described by the Heisenberg–Dirac–Van Vleck (HDVV) spin Hamiltonian, Equation (1):

190

$$191 \quad H = -J \times S_{\text{Ln}} \times S_{\text{Ln}} \quad (1)$$

192

193 But the spin Hamiltonian can only be used for GdIII, since it has no orbital contribution, and therefore,  
194 no spin-orbit coupling effect.[31] The other lanthanides(III) need much more complex models, based on  
195 explicit ligand field spin-orbit parameters.[32] Usually, such treatments are not carried out, due to their  
196 complexity, so the magnetic behaviour of 1–7 will be described, but the coupling constant  $J$  will only be  
197 calculated for the GdIII compound 3.

198 At room temperature, the  $M \times T$  values for compounds 1–7 are 3.14, 2.91, 16.13, 23.88, 25.95, 23.46  
199 and 4.93 cm<sup>3</sup> mol<sup>-1</sup> K, respectively. These data are in good agreement with the expected values for the  
200 corresponding two noninteracting LnIII–LnIII centres with ground states 4I<sub>9/2</sub> [NdIII], 7F<sub>0</sub> [EuIII],  
201 8S<sub>7/2</sub> [GdIII], 7F<sub>6</sub> [TbIII], 6H<sub>15/2</sub> [DyIII], 4I<sub>15/2</sub> [ErIII] and 2F<sub>7/2</sub> [YbIII].[9] Upon cooling, the  $M \times$   
202  $T$  values gradually decrease for 1, 2 and 7, which should be mainly attributed to the depopulation of  
203 their excited states. At 2.0 K, the  $M \times T$  value are 1.17, 0.03 and 2.18 cm<sup>3</sup> mol<sup>-1</sup> K for 1, 2 and 7,  
204 respectively, indicating an  $mJ = 0$  ground substate for EuIII ion (7F<sub>0</sub>).

205 For 3 and 4, the  $M \times T$  product remains almost constant down to ca. 30 K for Gd and ca. 45 K for Tb  
206 and then decreases to 13.41 cm<sup>3</sup> mol<sup>-1</sup> K (for 3) and to 13.50 cm<sup>3</sup> mol<sup>-1</sup> K (for 4) at 2.0 K, suggesting  
207 a moderately weak antiferromagnetic exchange interaction. In the case of compound 3, fitting of the  
208 experimental data was performed by means of Equation (1), by using the PHI computer program.[33]  
209 The best-fit parameters obtained were  $J = -0.04(1)$  cm<sup>-1</sup> and  $g = 2.01(1)$ , confirming the existence of  
210 antiferromagnetic coupling between the metal centres. The  $J$  value is comparable with those found for  
211 similar compounds.[12,34]

212 In complex 5, the values of the  $M \times T$  product increase slightly, up to a maximum value of 27.39 cm<sup>3</sup>  
213 mol<sup>-1</sup> K at 35 K. This increase of  $M \times T$  might indicate moderate intramolecular ferromagnetic  
214 exchange interactions. Below this temperature, the  $M \times T$  values decrease continuously down to 15.81  
215 cm<sup>3</sup> mol<sup>-1</sup> K at 2.0 K, which can be attributed mainly to the depopulation of the DyIII  $mJ$  sublevels of  
216 the ground  $J$  state.

217 The  $M \times T$  values for compound 6 continuously decrease when cooling, reaching a plateau at 13 K  
218 and then drops down to 16.26 cm<sup>3</sup> mol<sup>-1</sup> K at 2 K.

219

## 220 **Ac Magnetic Susceptibility Study**

221 Ac magnetic susceptibility measurements were performed on compounds 1, 4, 5, 6 and 7. Under zero dc  
222 magnetic fields, no maxima for the in-phase ( $M'$ ) and/or out-of-phase ( $M''$ ) susceptibility  
223 components were observed, probably due to an important quantum tunnelling of the magnetization  
224 (QTM) process present in these systems.[35] To suppress the QTM relaxation process, an optimal  
225 external dc field of 0.15 T for 1, 0.1 T for 5 and 6 and 0.2 T for 7 was applied and slow relaxation of the  
226 magnetization was then revealed (Figure 6).[24] In the case of compound 4, even under applied dc fields  
227 up to 0.4 T, no frequency or thermal dependencies of  $M'$  and/or  $M''$  are observed.

228 Representation of  $M''$ , measured at different frequencies (1–1488 Hz) of the  $0.4 \times 10^{-3}$  T ac field, at  
229 temperatures between 1.8 and 8.0 K, of compound 1 (Figure 6a) shows maxima of the out-of-phase

230 component. The ac curves recorded between 1.8 and 3.6 K have been well-fitted under the generalized  
 231 Debye model, as is shown in the corresponding Cole–Cole plot (Figure S1). The extracted relaxation  
 232 parameters are collected in Table S4, showing  $\alpha$  values close to zero, which means there is a narrow  
 233 distribution of the relaxation times and indicates that a single relaxation is mainly involved in the  
 234 system's relaxation process.[36] Thermal dependency of the relaxation time of the magnetization ( $\tau$ ) in  
 235 compound 1 (Figure 6b) reveals that at the highest temperatures, it follows an Orbach relaxation  
 236 process. The pre-exponential factor ( $\tau_0$ ) and the effective energy barrier ( $E_a$ ) between the two ground  
 237 magnetic states can then be extracted from Arrhenius  $\{\tau = \tau_0 \times \exp[E_a/(k_B \times T)]\}$ , giving values of  $2.8$   
 238  $\times 10^{-6}$  s and  $10.16$  cm $^{-1}$ , respectively, consistent with the values obtained for similar  
 239 compounds.[37,38] At low temperatures, the rate of  $\tau$  derives from linearity, probably due to the  
 240 presence of other relaxation mechanisms, such as Raman or direct processes.[39] The full temperature  
 241 range can be fitted with the following model, Equation (2):  $\tau^{-1} = \tau_0^{-1} \times \exp[-E_a/(k_B \times T)] + C \times T^n + A \times T$

$$\tau^{-1} = \tau_0^{-1} \times \exp[-E_a/(k_B \times T)] + C \times T^n + A \times T \quad (2)$$

245 where the first term represents an Orbach process, and the second and third ones represent Raman and  
 246 direct relaxation processes, respectively. The  $n$  parameter was fixed at 9, the usual value for Kramers  
 247 ions.[40] The best fit values obtained are:  $\tau_0 = 7.4 \times 10^{-6}$  s,  $E_a = 9.5$  cm $^{-1}$ ,  $C = 0.02$  s $^{-1}$  K $^{-9}$  and  $A =$   
 248  $265.21$  s $^{-1}$  K $^{-1}$ .

249 The ac susceptibility measurements were carried out on compound 5, with a  $0.4 \times 10^{-3}$  T ac field  
 250 oscillating at six frequencies between 10 and 1488 Hz, in the temperature range of 1.8–6.0 K. The out-  
 251 of-phase signal  $M''$  versus  $\nu$  and  $M''$  versus  $T$  plots are represented in Figure 6c and Figure S2,  
 252 respectively, showing the thermal and frequency dependencies of  $M''$ , but without net maxima in the  
 253 range of the measured temperatures, which indicates that the QTM is not completely suppressed. Then,  
 254 the energy barrier and relaxation time cannot be extracted from the Arrhenius equation; therefore, these  
 255 parameters were obtained using the Debye model and the equation  $\ln(M''/M') = \ln(\omega \times \tau_0) + E_a/(k_B$   
 256  $\times T)$ , [41,35] giving, from the best fit [see Figure S9 (right) and Table S5], a value of  $E_a = 6.6$  cm $^{-1}$  and  
 257 an exponential factor  $\tau_0 = 8.2 \times 10^{-7}$  s, in good agreement with similar DyIII dinuclear complexes.[41]  
 258 In complexes 6 and 7, frequency-dependent peaks for the  $M''$  component can be observed (Figure 6e  
 259 and g) under a  $0.4 \times 10^{-3}$  T ac field oscillating at frequencies between 1–1488 Hz, in the temperature  
 260 range of 1.8–4.5 K. Frequency dependences of both  $M'$  and  $M''$  below 3.3 K for 6 and 3.0 K for 7  
 261 were analyzed, again using the generalized Debye model.[42] The correspondent Cole–Cole plots  
 262 depicted in Figures S3 and S4 show almost semicircular shapes in all cases, with estimated  $\alpha$  values  
 263 close to zero. The best-fit relaxation parameters are collected in Tables S6 and S7.

264 The provided relaxation times ( $\tau$ ) allow the representations of  $\ln(\tau)$  versus  $T^{-1}$  (Figure 6f and h),  
 265 showing that at temperatures above 2.7 K for 6 and 2.2 K for 7, the rate of  $\tau$  follows the Arrhenius law,  
 266 giving energy barriers of 7.1 and 5.0 cm $^{-1}$  for 6 and 7, respectively, and pre-exponential factors ( $\tau_0$ ) of



267  $2.9 \times 10^{-6}$  s (6) and  $9.3 \times 10^{-6}$  s (7). These given values are in good agreement with other similar ErIII  
268 and YbIII field-induced SMM compounds found in the literature.[35,37,43] The deviation of the lineal  
269 dependence of  $\tau$  at low temperatures, for both compounds, suggests the combination of various  
270 relaxation processes of the magnetization. Data of compound 6 can be fitted using Equation (3),[44]  
271 affording values of  $C = 0.08 \text{ s}^{-1} \text{ K}^{-9}$  and  $A = 1296 \text{ s}^{-1} \text{ K}^{-1}$ .

$$\tau^{-1} = C \times T^9 + A \times T \quad (3)$$

275 For compound 7, the relaxation times at low temperatures can be modelled considering both Orbach and  
276 Raman processes; Equation (4):

$$\tau^{-1} = \tau_0^{-1} \exp[-E_a/(k_B \times T)] + C \times T^9 \quad (4)$$

280 The obtained values are  $\tau_0 = 2.2 \times 10^{-5}$  s,  $E_a = 3.7 \text{ cm}^{-1}$  and  $C = 0.09 \text{ s}^{-1} \text{ K}^{-9}$ .

## 282 Photoluminescence Properties

283 Excitation and emission spectra of complexes 1, 2 and 4–7, recorded in the solid state, at room  
284 temperature, are shown in Figure 7 and Figures S6–S9. The excitation spectra show broad bands around  
285 350 nm, corresponding to the  $\pi$ – $\pi^*$  transitions of the organic ligands in all of the compounds (see Figure  
286 S5). Since the GdIII ion has higher excited electronic states that cannot be sensitized with conventional  
287 organic ligands, the luminescence study of compound (3) reveals internal ligand transitions (Figure S7).  
288 In the cases of the NdIII, EuIII, TbIII and ErIII compounds, it is possible to observe the 4f–4f metal  
289 absorption peaks above 375 nm.

290 Luminescence measurements of complex 1 (Figure 7) exhibit emissions of the characteristic f–f  
291 transitions of the NdIII ion in the NIR region. The emission spectra monitored at  $\lambda_{\text{exc}} = 350$  nm presents  
292 a broad band centred at 416 nm, attributed to residual ligand emission in the visible region and the metal  
293 f–f transitions  $4F3/2 \rightarrow 4I9/2$  at 902 nm and  $4F3/2 \rightarrow 4I11/2$  at 1061 nm in the NIR region. From the  
294 excitation spectra recorded at the most intense NdIII transition  $4F3/2 \rightarrow 4I11/2$  ( $\lambda_{\text{em}} = 1061$  nm), the  
295 highest absorption band corresponds to ligand-centred transitions from 300 nm to 380 nm. Also, bands  
296 from direct NdIII absorption are revealed at 424 nm for  $2P1/2 \leftarrow 4I9/2$ , at 468 nm for  $2K15/2 + 4G11/2$   
297  $\leftarrow 4I9/2$ , at 513 nm for  $2K13/2 + 4G9/2 \leftarrow 4I9/2$ , at 526 nm for  $4G7/2 \leftarrow 4I9/2$ , at 588 nm for  $2G7/2 \leftarrow$   
298  $4I9/2$  and at 743 nm for  $4F7/2 + 4S3/2 \leftarrow 4I9/2$  transitions. The intensities ratios between the ligand and  
299 metal absorption peaks demonstrate that the NdIII emission in compound 1 is much more efficient when  
300 it is excited in the ligand absorption range than at direct metal absorption wavelengths. However, the  
301 ligand emission bands present in the emission spectra recorded in the ligand absorption range reveal that  
302 the energy transfer is not complete.

303 The excitation spectra of compound 2 (Figure S6), recorded by monitoring the 5D0 → 7F2 transition of  
304 EuIII at  $\lambda_{em} = 615$  nm, shows broad ligand-centred  $\pi^* \leftarrow \pi$  transition bands from 316 to 370 nm. At  
305 lower energies, it is possible to distinguish the f–f transitions of the metal displaying considerably less  
306 intensity with respect to the ligand transitions, suggesting that the emission is supported by the antenna  
307 effect involving the organic fragment. The emission spectra of 2 (Figure S6), recorded under excitation  
308 at 350 nm, results in the characteristic luminescence of 5D0 → 7FJ (J = 0–4) transitions from the EuIII  
309 ion. Specific assignments are as follow: 5D0 → 7F0 (579 nm), 5D0 → 7F1 (592 nm), 5D0 → 7F2 (615  
310 nm), 5D0 → 7F3 (649 nm) and 5D0 → 7F4 (697 nm). The red emission light spectrum is dominated by  
311 the hypersensitive 5D0 → 7F2 transition, which shows signs of splitting. This fact indicates that the  
312 chemical environment around the EuIII ion does not present an inversion centre.[9,45] Moreover, the  
313 presence of only one sharp peak corresponding to the 5D0 → 7F0 transition reveals that all of the EuIII  
314 centres present the same chemical environment, and thus, the same crystal field, which is in good  
315 agreement with the X-ray diffraction crystallographic data. The absence of ligand-centred emission  
316 bands suggests that the antenna effect in this complex is efficient.

317 In the case of the TbIII compound 4, the excitation spectrum recorded at  $\lambda_{em} = 545$  nm, corresponding  
318 to the emission band of 5D4 → 7F5 transition, shows the ligand absorption bands and weak peaks  
319 coming from some f–f transitions from the metal (Figure S8). Complex 4 shows the typical TbIII  
320 emission spectrum, containing the expected sequence of 5D4 → 7FJ (J = 3–6) transitions (Figure S8).  
321 Specific assignments are as follows: 5D4 → 7F6 (489 nm), 5D4 → 7F5 (545 nm), 5D4 → 7F4 (585  
322 nm), 5D4 → 7F3 (620 nm) and 5D4 → 7F2 (646 nm). The spectrum is dominated by the 5D4 → 7F5  
323 transition, which gives an intense green luminescence output for the solid sample. As for the EuIII  
324 complex, the energy transfer between the ligand and the metal in this complex seems to be efficient, due  
325 to the absence of ligand emission bands.

326 The ligand  $\pi^* \leftarrow \pi$  transition dominates the excitation spectrum of compound 5 (Figure S9) measured at  
327 the 4F9/2 → 6H13/2 transition ( $\lambda_{em} = 572$  nm). Then, the antenna effect enhances the f–f emission  
328 peaks of the DyIII at 478 and 572 nm under excitation at 350 nm, which are assigned to the 4F9/2 →  
329 6HJ (J = 15/2, 13/2 transitions, respectively) (Figure S9). The spectrum is dominated by the 4F9/2 →  
330 6H15/2 transition, which gives an intense blue luminescence output for the solid sample.

331 Compound 6 shows intense ligand-centred emission at around 416 nm (Figure 7). In the NIR range, this  
332 compound also exhibits the characteristic emission peak corresponding to the transition 4I13/2 → 4I15/2  
333 at 1531 nm. In the excitation spectra (Figure 7), there are bands corresponding to through-ligand  
334 excitation (EL) and others are assigned to through-metal excitation (ELn) at lower energies. The  
335 EL/ELn intensity ratio in the NIR emitting complex 6 is lower than those from the EuIII and TbIII  
336 complexes. This fact is because the f–f transitions for ErIII are less Laporte forbidden, so the direct  
337 lanthanide excitation is more significant in this compound.[46]

338 For complex 7, ligand-centred emission is also observed around 416 nm (Figure 7), and at 993 nm, the  
339 YbIII 2F5/2 → 2F7/2 transition is revealed. The excitation spectrum of 7 ( $\lambda_{em} = 993$  nm) shows only

340 the absorption of the ligand; as the YbIII ion does not present electronic levels in the UV/Vis region, the  
341 sensitization of the LnIII emission could be only caused by the ligand antenna effect (Figure 7).[47]

342

343

344

345 **CONCLUSIONS**

346

347 Here, we have presented structural, magnetic and luminescence studies of a family of homo-dinuclear  
348 lanthanide compounds based on the use of 2-FBz and phen ligands. We have used a straightforward  
349 room-temperature synthetic procedure and we have successfully obtained six new dinuclear lanthanide  
350 compounds. From a structural point of view, the diminution of the ionic radius along the 4f row is  
351 translated into two different structures being obtained. Following the same synthetic method; for large  
352 ionic radii, compounds 1–6, with the general formula  $[\text{Ln}_2(2\text{-FBz})_4(\text{NO}_3)_2(\text{phen})_2]$  ( $\text{Ln} = \text{Nd}, \text{Eu}, \text{Gd},$   
353  $\text{Tb}, \text{Dy}$  and  $\text{Er}$ ) present nine-coordinate environment metal ions. Meanwhile, in compound  $[\text{Yb}_2(2\text{-}$   
354  $\text{FBz})_6(\text{phen})_2]$  (7), with YbIII having the smallest ionic radius of the series, the coordination number  
355 falls to eight, completely changing the ion environment.

356 All of the compounds discussed in this work, except for the GdIII one, display the corresponding f–f  
357 emission luminescence, due to ligand absorption followed by energy transfer to the metal, direct metal  
358 absorption or a combination of the two procedures. The EuIII (2), TbIII (4) and DyIII (5) products emit  
359 in the visible range, and the NdIII (1), ErIII (6) and YbIII (7) compounds emit in the NIR region.

360 The fitting of the  $M \times T$  versus  $T$  curve of the GdIII (3) compound by the Heisenberg–Dirac–Van  
361 Vleck (HDVV) spin Hamiltonian reveals a weak antiferromagnetic interaction between the two LnIII  
362 within the dinuclear unit. Additionally, dynamic magnetic measurements reveal field-induced SMM  
363 character for compounds 1 and 5–7.

364 Thus, compounds 1, 5, 6 and 7 present both field-induced SMM and luminescence properties and are  
365 considered to be multifunctional complexes with potential biomedical applications.

366

367 **EXPERIMENTAL SECTION**

368  
369 **Starting Materials:** Ln(NO<sub>3</sub>)<sub>3</sub>·6H<sub>2</sub>O salts, 2-fluorobenzoic acid and 1,10-phenanthroline (Aldrich)  
370 were used as received, without further purification.

371  
372 **General Syntheses**

373 **Complexes 1–7** were obtained on the basis of a previously proposed synthetic approach.[48] The  
374 preparation of all compounds was achieved by the reaction of 2-fluorobenzoic acid (H-2-FBz, 3 mmol)  
375 and 1,10-phenanthroline (phen, 0.6 mmol) dissolved in EtOH/H<sub>2</sub>O (20 mL, v/v = 50:50) with a solution  
376 of the corresponding Ln(NO<sub>3</sub>)<sub>3</sub>·6H<sub>2</sub>O salt (0.5 mmol) [Ln = Nd (1), Eu (2), Ln = Gd (3), Ln = Tb (4),  
377 Ln = Dy (5), Ln = Er (6), Ln = Yb (7)] in EtOH (10 mL). The mixture was stirred for 1 h at room  
378 temperature. Good-quality crystals, suitable for X-ray diffraction analysis of compounds 1–7, were  
379 obtained after 5–10 days of slow evaporation.

380  
381 **Compound 1:** C<sub>52</sub>H<sub>32</sub>F<sub>4</sub>N<sub>6</sub>Nd<sub>2</sub>O<sub>14</sub> (1329.32): calcd. C 46.98, H 2.43, N 6.32; found C 46.5, H 2.5, N  
382 6.8. Selected IR bands (KBr pellet):  $\tilde{\nu}$  = 1613 (vs), 1592 (s), 1395 (vs), 1384 (vs), 1305 (s) cm<sup>-1</sup>.

383  
384 **Compound 2:** C<sub>52</sub>H<sub>32</sub>Eu<sub>2</sub>F<sub>4</sub>N<sub>6</sub>O<sub>14</sub> (1344.76): calcd. C 46.44, H 2.40, N 6.25; found C 46.8, H 2.4, N  
385 6.5. Selected IR bands (KBr pellet):  $\tilde{\nu}$  = 1618 (vs), 1565 (s), 1397 (vs), 1385 (vs), 1305 (s) cm<sup>-1</sup>.

386  
387 **Compound 3:** C<sub>52</sub>H<sub>32</sub>Gd<sub>2</sub>F<sub>4</sub>N<sub>6</sub>O<sub>14</sub> (1355.34): for C 46.08, H 2.38, N 6.20; found C 45.3, H 2.3, N  
388 6.1. Selected IR bands (KBr pellet):  $\tilde{\nu}$  = 1615 (vs), 1563 (s), 1395 (vs), 1384 (vs), 1306 (s) cm<sup>-1</sup>.

389  
390 **Compound 4:** C<sub>52</sub>H<sub>32</sub>F<sub>4</sub>N<sub>6</sub>O<sub>14</sub>Tb<sub>2</sub> (1358.67): calcd. C 45.97, H 2.37, N 6.18; found C 46.4, H 2.4, N  
391 6.3. Selected IR bands (KBr pellet):  $\tilde{\nu}$  = 1615 (vs), 1563 (s), 1399 (vs), 1384 (vs), 1306 (vs) cm<sup>-1</sup>.

392  
393 **Compound 5:** C<sub>52</sub>H<sub>32</sub>Dy<sub>2</sub>F<sub>4</sub>N<sub>6</sub>O<sub>14</sub> (1365.84): calcd. C 45.73, H 2.36, N 6.15; found C 47.5, H 2.5, N  
394 5.5. Selected IR bands (KBr pellet):  $\tilde{\nu}$  = 1618 (vs), 1566 (w), 1395 (s), 1384 (vs), 1306 (s) cm<sup>-1</sup>.

395  
396 **Compound 6:** C<sub>52</sub>H<sub>32</sub>Er<sub>2</sub>F<sub>4</sub>N<sub>6</sub>O<sub>14</sub> (1375.36): calcd. C 45.41, H 2.34, N 6.11; found C 45.6, H 2.5, N  
397 6.0. Selected IR bands (KBr pellet):  $\tilde{\nu}$  = 1618 (vs), 1569 (s), 1395 (s), 1384 (s), 1306 (s) cm<sup>-1</sup>.

398  
399 **Compound 7:** C<sub>66</sub>H<sub>40</sub>F<sub>6</sub>N<sub>4</sub>O<sub>12</sub>Yb<sub>2</sub> (1541.10): calcd. C 51.44, H 2.62, N 3.63; found C 51.0, H 2.8, N  
400 3.8. Selected IR bands (KBr pellet):  $\tilde{\nu}$  = 1612 (vs), 1594 (s), 1415 (vs), 1288 (s) cm<sup>-1</sup>.

401  
402 **Spectral and Magnetic Measurements:** The elemental analyses of the compounds were performed at  
403 the Serveis Científics i Tecnològics of the Universitat de Barcelona. Infrared spectra (4000–400 cm<sup>-1</sup>)

404 were recorded from KBr pellets with a Perkin–Elmer 380-B spectrophotometer. Solid-state fluorescence  
405 spectra were recorded with a Horiba Jobin Yvon SPEX Nanolog fluorescence spectrophotometer  
406 at room temperature. Magnetic measurements were performed on solid polycrystalline samples in a  
407 Quantum Design MPMS-XL SQUID magnetometer at the Magnetic Measurements Unit of the  
408 Universitat de Barcelona. Pascal's constants were used to estimate the diamagnetic corrections, which  
409 were subtracted from the experimental susceptibilities to give the corrected molar  
410 magnetic susceptibilities.

411

412 **X-ray Diffraction Crystallography:** Crystals of 1–3 and 5–7 were mounted in air in a D8VENTURE  
413 (Bruker) diffractometer with CMOS detector. The crystallographic data, conditions retained for the  
414 intensity data collection and some features of the structural refinements are listed in Table 2. All of the  
415 structures were refined by the least-squares method. Intensities were collected with a multilayer  
416 monochromated Mo-K $\alpha$  radiation. Lorentz polarization and absorption corrections were made in all of  
417 the samples. The structures were solved by direct methods, using the SHELXS-97 computer  
418 program[49] and were refined by the full-matrix least-squares method, using the SHELXL-2014  
419 computer program.[50] The non-hydrogen atoms were located in successive difference Fourier  
420 syntheses and were refined with anisotropic thermal parameters on F2. For hydrogen atoms, isotropic  
421 temperature factors have been assigned as being 1.2 or 1.5 times that of the respective parent.

422

423 CCDC 1822670 (for 1), 1577089 (for 2), 1577090 (for 3), 1577091 (for 5), 1577734 (for 6), and  
424 1577735 (for 7) contain the supplementary crystallographic data for this paper. These data can be  
425 obtained free of charge from The Cambridge Crystallographic Data Centre.

426

427 **Supporting Information** (see footnote on the first page of this  
428 article): Tables S1–S7, Figures S1–S9 and crystallographic data in CIF  
429 or other electronic format.

430

431 **ACKNOWLEDGEMENTS**

432

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435

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- 524

525 **Legends to figures**

526

527 **Scheme 1** Representation of the different coordination modes of the 9-AC ligand.

528

529 **Figure. 1** (a) Partially labelled plot of compound 2. H atoms are omitted for clarity. Colour code: grey =  
530 C, blue = N, red = O, yellow = F and pink = Eu. (b) Coordination polyhedron of the EuIII ions in  
531 compound 2..

532

533 **Figure.2** Supramolecular arrangement representation in compound 2 on the (111) plane. In blue:  $\pi$ -  
534 stacking interactions; in green: C–H $\cdots$ F intramolecular  
535 hydrogen bonds.

536

537 **Figure.3.** (a) Partially labelled plot of compound 7. H atoms are omitted for clarity. Colour code: grey =  
538 C, blue = N, red = oxygen, yellow = fluorine and green = Yb. (b) Coordination polyhedron of the YbIII  
539 ions in compound 7.

540

541 **Figure.4** Representation of the supramolecular arrangement of compound 7 on the (111) plane. In blue:  
542  $\pi$ -stacking interactions; in green: C–H $\cdots$ F hydrogen bonds.

543

544 **Figure.5**  $M \times T$  versus  $T$  plots for compounds 1–7. The solid blue line corresponds to the best fit for  
545 complex 3 (see text).

546

547 **Figure 6** Representation of the frequency dependence of the out-of-phase component of the ac  
548 susceptibility under 0.15 T dc field for 1 (a), under 0.1 T dc field for 5 (c) and 6 (e) and under a 0.2 T dc  
549 field for 7 (g). Magnetization relaxation time,  $\ln(\tau)$  versus  $T^{-1}$  for 1 (b), for 6 (f) and for 7 (h). The red  
550 and blue lines represent the fitting using the Arrhenius equation and Orbach, Raman and/or direct  
551 relaxation processes, respectively (see text). (d) Representation of the natural logarithm of  $M''/M'$   
552 versus  $T^{-1}$  for 5. The red line is fitted using the following equation:  $\ln(M''/M') = \ln(\omega \times \tau_0) +$   
553  $E_a/(k_B \times T)$ .

554

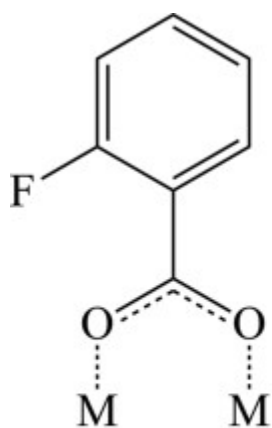
555 **Figure.7** Solid-state excitation (red) and emission (black) spectra of compounds 1, 6 and 7 at r. t.

556

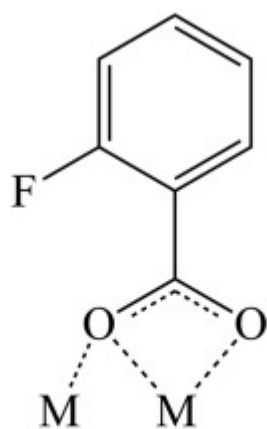
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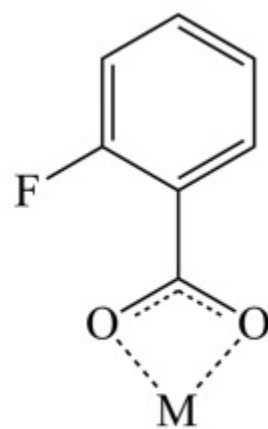
**SCHEME 1**



a) *syn, syn* bridge



b) chelating-bridging

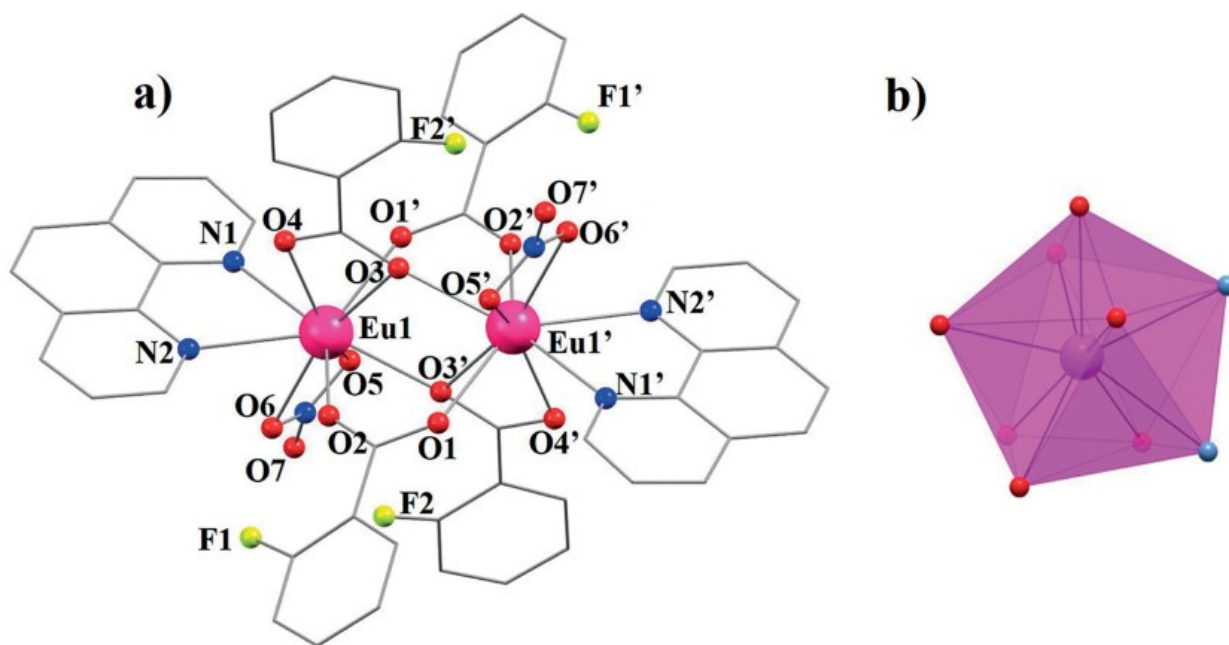


c) chelating

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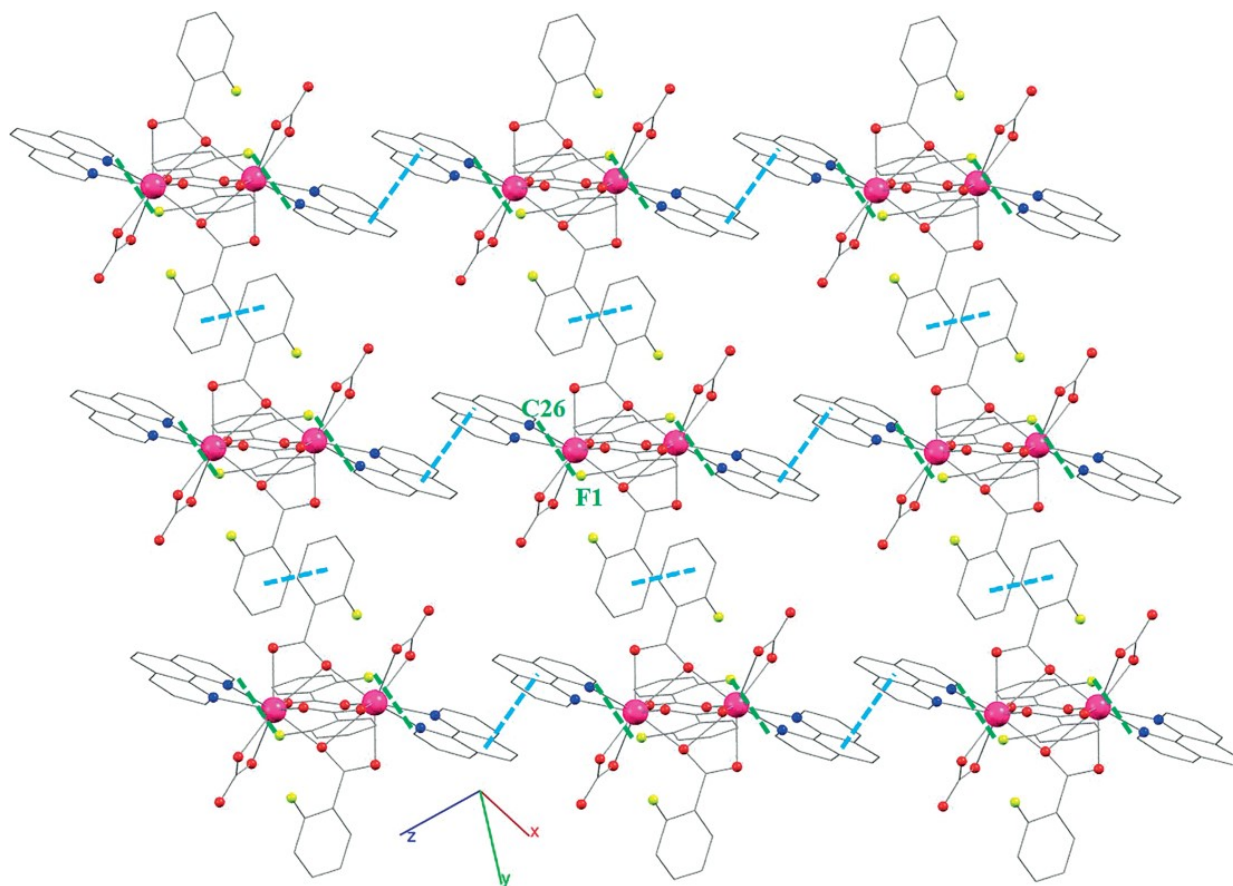
FIGURE 1



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FIGURE 2



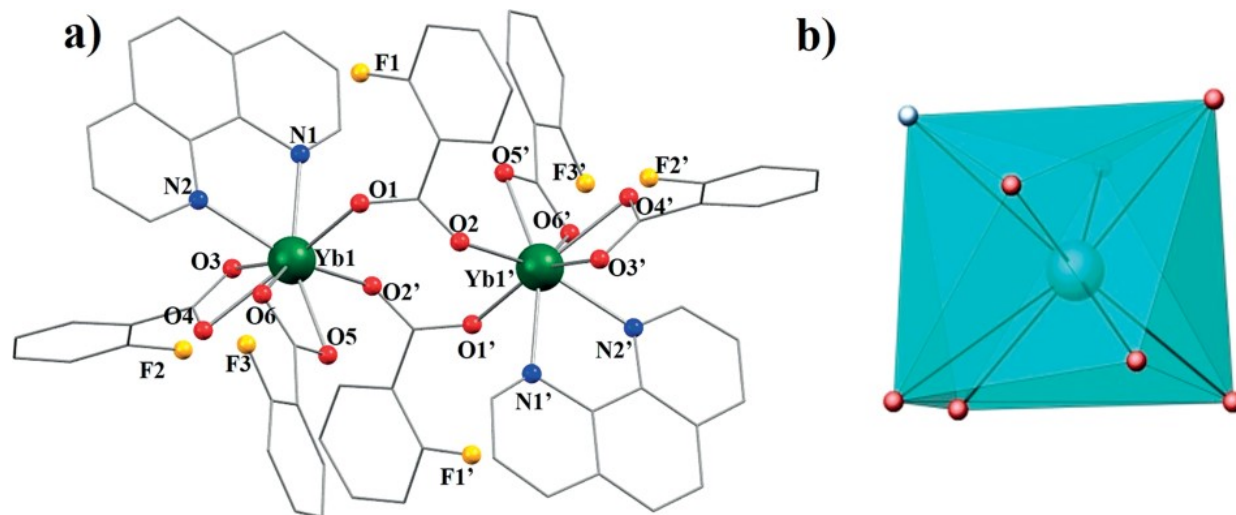
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FIGURE 3



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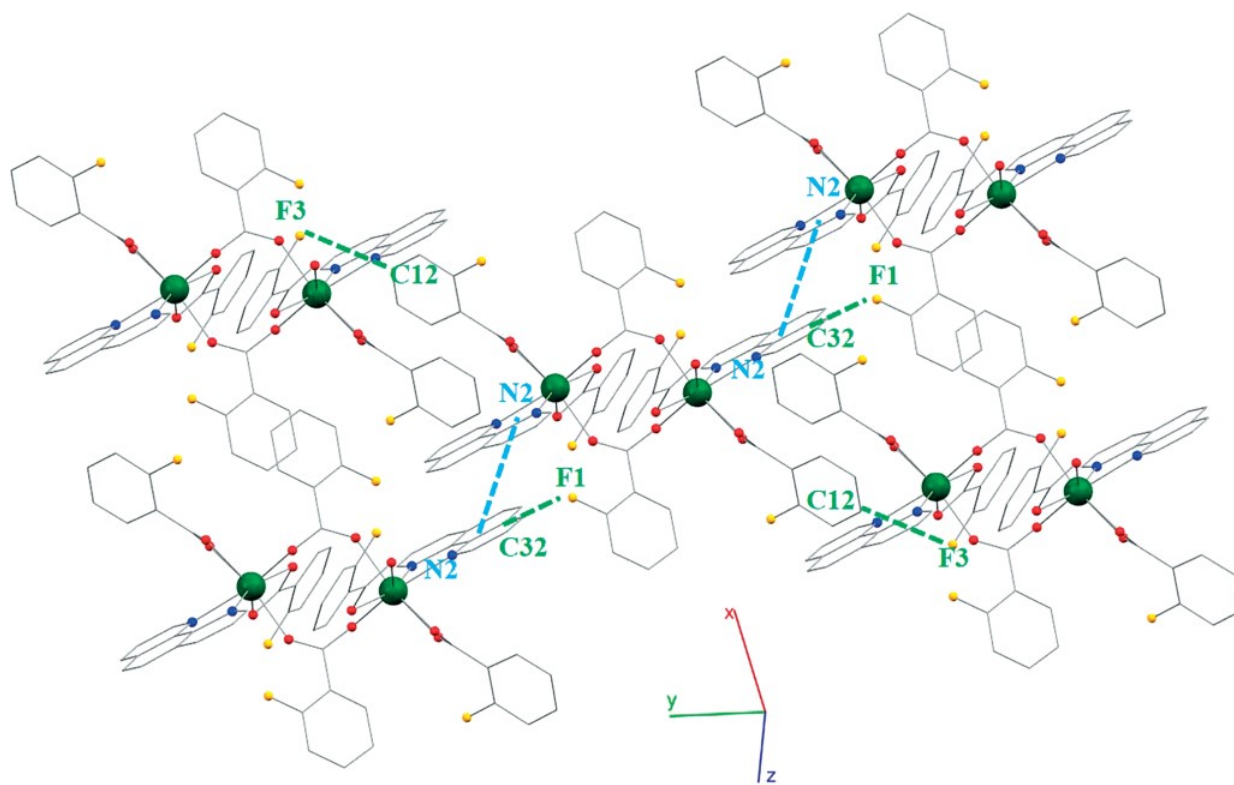


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FIGURE 4

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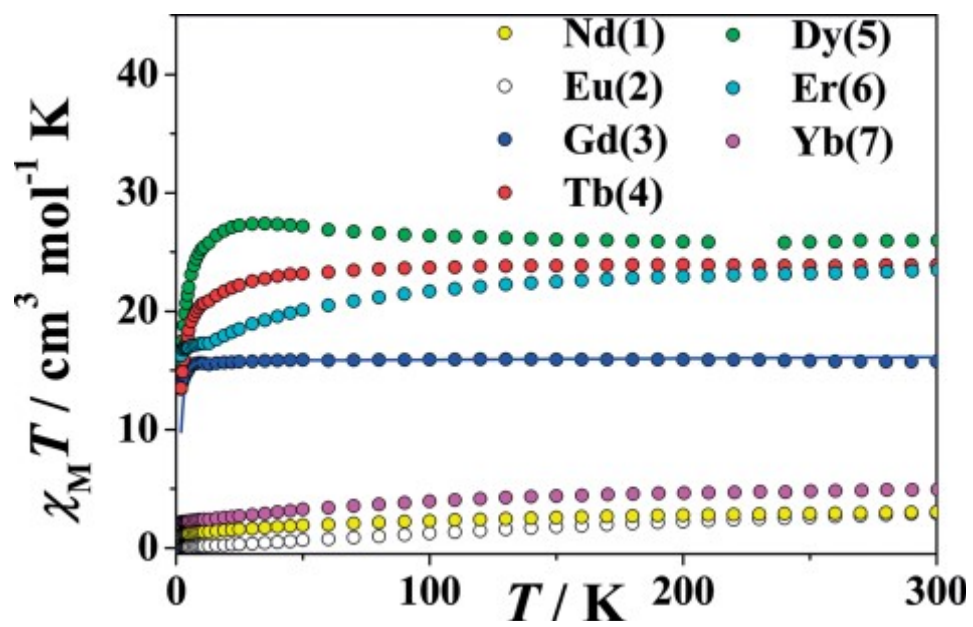
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FIGURE 5

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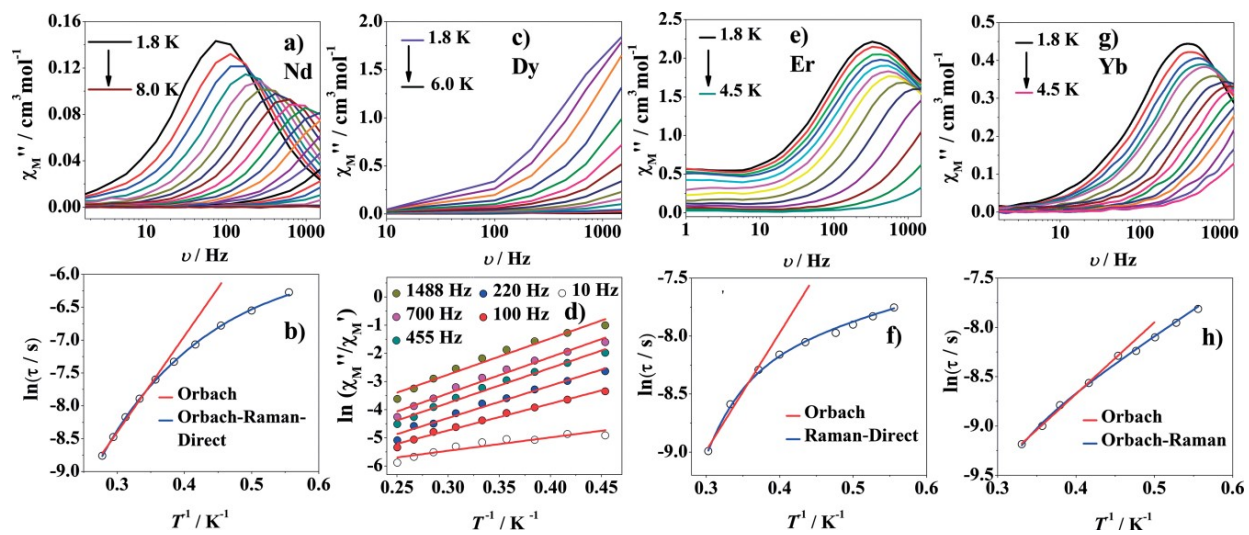
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FIGURE 6

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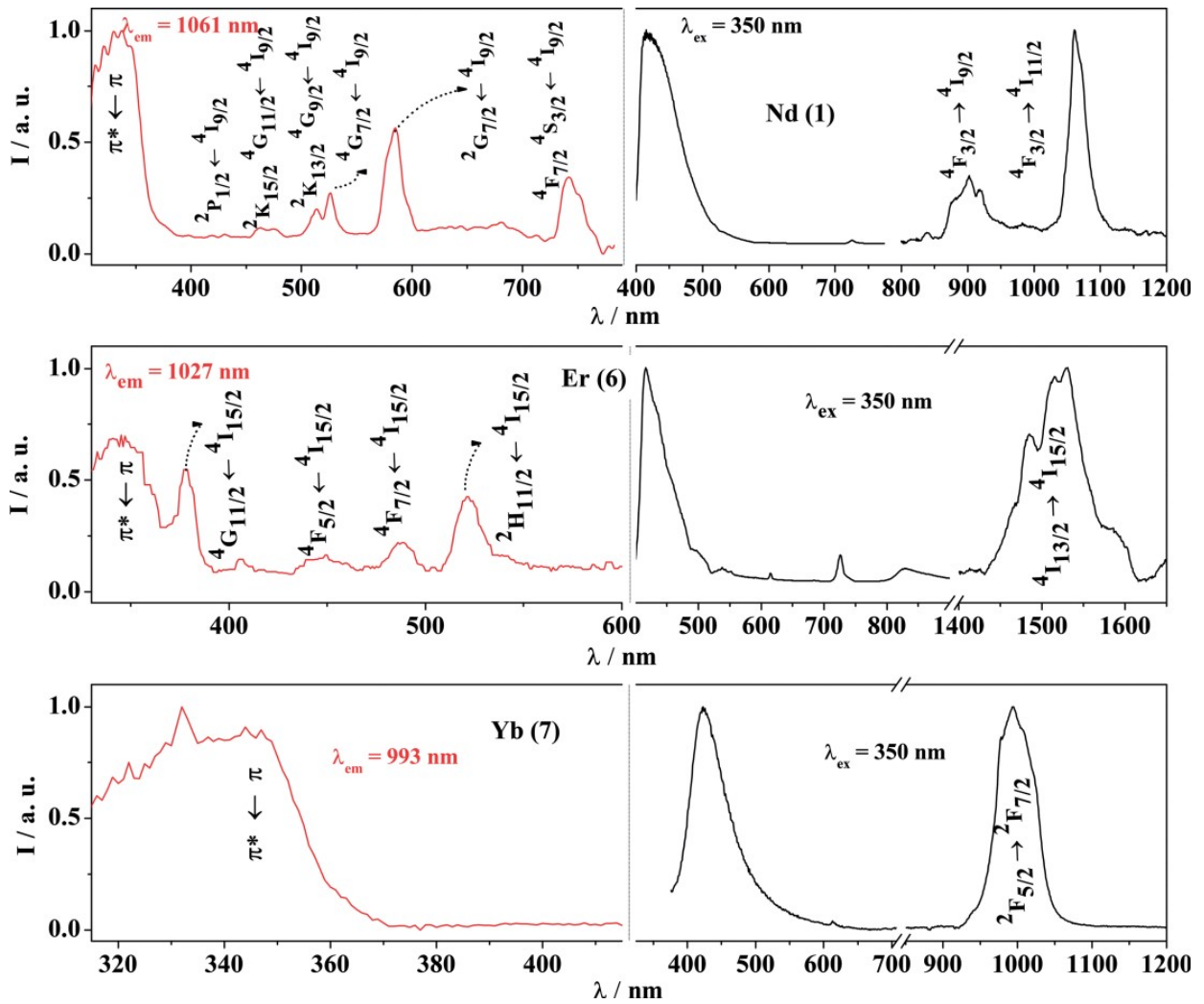
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FIGURE 7

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600 **Table 1.** Selected bond lengths [ $\text{\AA}$ ] for compounds 1–7.

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	1	2	3	4 <sup>(26)</sup>	5	6	7
<b>Ln1–O1</b>	2.3916(14)	2.3698(1)	2.2573(1)	2.337(2)	2.3283(2)	2.3053(12)	2.262(2)
<b>Ln1–O3</b>	2.4140(15)	2.3705(1)	2.2588(1)	2.340(2)	2.3265(2)	2.3008(13)	2.409(2)
<b>Ln1–O4</b>	2.4597(15)	2.4066(1)	2.2932(1)	2.374(2)	2.3567(2)	2.3301(13)	2.376(2)
<b>Ln1–O5</b>	2.5671(15)	2.5329(1)	2.5259(1)	2.514(4)	2.5056(2)	2.4923(13)	2.339(2)
<b>Ln1–O6</b>	2.5089(16)	2.4604(1)	2.4478(1)	2.423(4)	2.4152(2)	2.3853 (12)	2.380(2)
<b>Ln1–N1</b>	2.6249(16)	2.5811(1)	2.5711(1)	2.572(4)	2.5457(2)	2.5208(13)	2.459(2)
<b>Ln1–N2</b>	2.6519(17)	2.6170(1)	2.6094(1)	2.608(5)	2.5862(2)	2.5659(15)	2.385(6)
<b>Ln1–O2'</b>	2.4151(14)	2.3521(1)	2.2398(1)	2.327(2)	2.3120(2)	2.2888(12)	2.206(2)
<b>Ln1–O3'</b>	2.7052(14)	2.6989(1)	2.6971(1)	2.751(2)	2.7326(2)	2.7432(12)	–
<b>Ln1...Ln1'</b>	4.0015(4)	3.9605(1)	3.9517(3)	3.9713(8)	3.9497(4)	3.9294(3)	5.1519(7)
<b>Symmetry</b>	1 - x, 1 - y, 1 - z	1 - x, 2 - y, -z	-x, 2 - y, 1 - z	-x, -y, -z	1 - x, 1 - y, 1 - z	1 - x, 1 - y, 1 - z	1 - x, 1 - y, 2 - z

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604 **Table 2** Crystal data and collection details for the X-ray diffraction structure of complexes 1–3 and 5–7.

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	1	2	3	5	6	7
Formula	C <sub>52</sub> H <sub>32</sub> F <sub>4</sub> N <sub>6</sub> Nd <sub>2</sub> O <sub>14</sub>	C <sub>52</sub> H <sub>32</sub> Eu <sub>2</sub> F <sub>4</sub> N <sub>6</sub> O <sub>14</sub>	C <sub>52</sub> H <sub>32</sub> Gd <sub>2</sub> F <sub>4</sub> N <sub>6</sub> O <sub>14</sub>	C <sub>52</sub> H <sub>32</sub> Dy <sub>2</sub> F <sub>4</sub> N <sub>6</sub> O <sub>14</sub>	C <sub>52</sub> H <sub>32</sub> Er <sub>2</sub> F <sub>4</sub> N <sub>6</sub> O <sub>14</sub>	C <sub>66</sub> H <sub>40</sub> F <sub>6</sub> N <sub>4</sub> O <sub>12</sub> Yb <sub>2</sub>
Formula mass [g mol <sup>-1</sup> ]	1329.22	1344.76	1355.34	1365.84	1375.36	1541.10
System	triclinic	triclinic	triclinic	triclinic	triclinic	triclinic
Space group	P1	P1	P1	P1	P1	P1
a [Å]	10.6739(7)	10.6529(2)	10.6487(3)	10.6699(7)	10.6433(4)	11.5251(15)
b [Å]	10.7928(7)	10.8226(3)	10.8222(3)	10.8859(7)	10.8789(4)	11.6743(11)
c [Å]	11.2992(7)	11.1880(3)	11.1560(3)	11.1024(7)	11.0468(4)	11.9243(11)
α [°]	83.232(2)	83.521(2)	83.599(2)	83.870(2)	83.902(1)	114.216(8)
β [°]	82.070(2)	81.661(2)	81.561(2)	81.282(2)	81.129(1)	99.451(5)
γ [°]	70.223(2)	70.001(2)	69.942(2)	69.645(2)	69.463(1)	92.998(5)
V [Å <sup>3</sup> ]	1209.82(14)	1196.63(5)	1192.10(6)	1193.03(13)	1181.58(8)	1430.5(2)
Z	1	1	1	1	1	1
T [K]	100(2)	100(2)	100(2)	100(2)	100(2)	100(2)
λ(Mo-Kα) [Å]	0.71073	0.71073	0.71073	0.71073	0.71073	0.71073
D <sub>calc</sub> [g cm <sup>-3</sup> ]	1.825	1.866	1.886	1.901	1.933	1.789
μ(Mo-Kα) [mm <sup>-1</sup> ]	2.215	2.691	2.849	3.202	3.622	3.339
R	0.0568 (7009)	0.0237 (6769)	0.0152 (13986)	0.0254 (7176)	0.0151 (6938)	0.0337 (7481)
wR <sub>2</sub>	0.0206 (7260)	0.0562 (7278)	0.0252 (15022)	0.0648 (7312)	0.0384 (7208)	0.0709 (8878)

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