

Comblike Poly(α -alkyl γ -glutamate)s: Computer Simulation Studies of an Intermediate Thermal Phase

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Monte Carlo (MC) simulations have been used to study the structure of an intermediate thermal phase of poly(α -octadecyl γ ,D-glutamate). This is a comblike poly(γ -peptide) able to adopt a biphasic structure that has been described as a layered arrangement of backbone helical rods immersed in a paraffinic pool of polymethylene side chains. Simulations were performed at two different temperatures (348 and 363 K), both of them above the melting point of the paraffinic phase, using the configurational bias MC algorithm. Results indicate that layers are constituted by a side-by-side packing of 17/5 helices. The organization of the interlayer paraffinic region is described in atomistic terms by examining the torsional angles and the end-to-end distances for the octadecyl side chains. Comparison with previously reported comblike poly(β -peptide)s revealed significant differences in the organization of the alkyl side chains.

Introduction

Helical comblike polypeptides are of interest because of their peculiar structure.¹ These polymers are constituted by two different structural units: (i) a rigid helical backbone, which is stabilized by intramolecular hydrogen bonds, and (ii) a flexible long linear alkyl side chain, the conformation of which mainly depends on the temperature. As a result, comblike polypeptides adopt a biphasic structure, which consists of a layered arrangement of backbone helical rods immersed in a paraffinic pool (Figure 1). This structure is highly sensitive to temperature, so phase transitions are usually induced by heating or cooling treatments. The main structural differences between these phases concern the organization of the alkyl side chains in the paraffinic region. In the most ordered phase, denoted A, the alkyl side chains are partially crystallized in a hexagonal lattice with the interior methylene units in trans conformation. Phase A converts into phase B upon heating above the transition temperature, T_1 . In phase B, the alkyl side chains are in a molten state even though the polypeptide main chains retain the helical conformation and the layered arrangement. A third phase C is observed at higher temperatures in certain cases, but the structure present in this phase is not well understood.

A notable effort was dedicated in the last 2 decades to investigate the structure of comblike poly(γ -alkyl α ,L-glutamate)s, that is, the alkyl esters of poly(α -glutamic acid), abbreviated PGALG- n (where n indicates the number of carbon atoms of the alkyl side chain).¹⁻³ Recently, we evidenced that helical comblike poly(β -peptide)s behave similarly to PGALG- n .⁴⁻⁸ More specifically, poly(α -alkyl β ,L-aspartate)s, abbreviated PAALA- n , were examined by

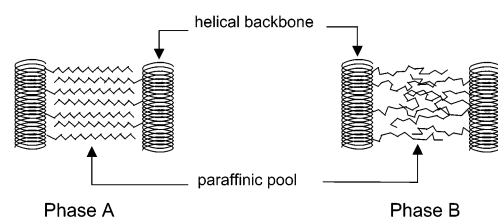
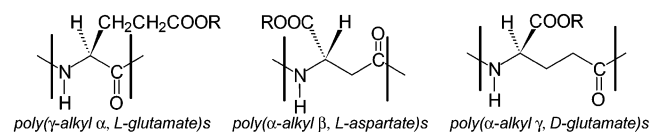


Figure 1. Schematic model illustrating the phases A and B of comblike polypeptides.

differential scanning calorimetry, NMR, and X-ray diffraction and shown to adopt the characteristic biphasic structure with the polypeptide chains in 13/4 helical conformation.⁴⁻⁶ Furthermore, an atomistic structural description was provided for comblike PAALA- n by using Monte Carlo (MC) simulations.⁵⁻⁸



On the other hand, the synthesis and structure of poly(α -alkyl γ -glutamate)s with short and long alkyl side chains, abbreviated PAAG- n , have been examined by us.⁹⁻¹² These poly(γ -peptide)s were prepared either by chemical synthesis or by derivatization of bacterially produced poly(γ -glutamic acid) with different D/L enantiomeric ratios. A recent analysis of the microstructure of these biosynthetic polymers by ¹³C NMR revealed that they are actually stereocopolymers made of enantiomerically homogeneous D and L blocks.¹⁰ Accordingly, optically pure polymers and racemic polymers exhibited practically the same structural behavior.

Two series of PAAG- n differing in the D/L enantiomeric ratio, 9:1 and 1:1, and with n ranging from 12 to 22 were synthesized and characterized by a variety of experimental techniques.^{9a} In a preceding paper, a detailed experimental

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68 study is carried out on the supramolecular structure adopted
 69 by these comblike poly(γ -glutamate)s with particular atten-
 70 tion paid to the phase transitions that take place by effect of
 71 temperature.^{9b} The observations were consistent with the
 72 occurrence of a biphasic layered structure, the layers being
 73 constituted by helices stabilized by intramolecular hydrogen
 74 bonds. Such results along with those previously obtained for
 75 poly(β -peptide)s establish that this type of supramolecular
 76 assemblies is not found only for α -helix but that it is shared
 77 by other nonconventional polypeptides provided that they
 78 can be arranged in a helical conformation. The investigation
 79 of PAAG-*n* is particularly interesting because of the bio-
 80 synthetic accessibility and potential biodegradability of their
 81 parent compound, poly(γ -glutamic acid) (PGGA). Accord-
 82 ingly, the structure of PAAG-*n* deserves a microscopic study
 83 using atomistic simulations.

84 The purpose of this work is to provide a detailed
 85 description of the atomistic structure of comblike PAAG-*n*
 86 using advanced MC methods. The study has been confined
 87 to the study of the structure of phase B because of (i) the
 88 difficulty of MC technique to reproduce the crystallized
 89 region of phase A, (ii) the almost total lack of experimental
 90 information on the structure of phase C, and (iii) the amount
 91 of experimental data available on phase B that, although
 92 insufficient to attain a detailed description of the structure
 93 adopted in this phase, is very useful for supporting the
 94 simulation analysis. Simulations have been performed con-
 95 sidering the enantiomerically pure poly(α -octadecyl γ ,*D*-
 96 glutamate), denoted PAADG-18, for consistency with our
 97 previous studies on PAALA-18.^{5,7,8} Furthermore, a detailed
 98 comparison between helical comblike poly(β -peptide)s and
 99 poly(γ -peptide)s has been made.

100 Model and Computational Methods

101 **Helical Conformation.** Precise experimental information
 102 required to define the helical backbone conformation of
 103 comblike polypeptides is not usually attained because of the
 104 constitutional complexity of these systems. For comblike
 105 poly(α -peptide)s and poly(β -peptide)s, the molecular con-
 106 formation was inferred from the X-ray data obtained for
 107 members with short alkyl side groups, which crystallize in
 108 a three-dimensional array.^{1,2,4} Accordingly, 18/5 and 13/4
 109 helical conformations were proposed for comblike PGALG-*n*
 110 and PAALA-*n*, respectively. In the present study, we have
 111 initially considered the helical conformations previously
 112 described for poly(γ -glutamic acid) and poly(γ -glutamate)s
 113 bearing short side chains. Experimental observations first
 114 carried out by Rydon,¹⁵ and later by other authors, evidenced
 115 that poly(γ ,*D*-glutamic acid) adopts a helical conforma-
 116 tion.^{16,17} Computer simulations revealed recently that such
 117 experimental data are compatible with a left-handed 17/5
 118 helix stabilized by intramolecular hydrogen bonds set
 119 between the amide groups *i* and *i* + 3.¹⁸ On the other hand,
 120 synthetic methyl and benzyl esters of poly(γ ,*L*-glutamic acid)
 121 were found to adopt a 5/2 helical conformation stabilized
 122 by intramolecular hydrogen bonds between the *i* and *i* + 2
 123 amide groups.^{11,12} Solution NMR experiments carried out on
 124 related oligo(γ ,*L*-amino acid)s detected the same conforma-

Table 1. Conformational Angles and Hydrogen-Bonding Parameters for the Helical Conformations Considered in This Work

	helix		
	5/2	17/5	37/10
dihedral angles ^a			
φ	-137.9	70.9	70.9
ξ_1	53.0	52.8	53.6
ξ_2	75.6	-171.0	-172.8
ψ	-143.8	159.9	164.0
ω	180.0	180.0	180.0
η	180.0	155.0	157.0
rise per residue ^b	2.02	1.50	1.50
helix sense	right-handed	left-handed	left-handed
H-bond type	intramolecular	intramolecular	intramolecular
H-bond scheme ^c	C=O(<i>i</i>)... H-N(<i>i</i> + 2)	N-H(<i>i</i>)... O=C(<i>i</i> + 3)	N-H(<i>i</i>)... O=C(<i>i</i> + 3)
$d(\text{H}\cdots\text{O})^b$	1.82	1.90	2.02
$\angle\text{N-H}\cdots\text{O}^a$	164.8	167.0	165.1
atoms per H-bond	14	19	19

^a In deg. ^b In Å. ^c The label indicated in parentheses corresponds to the number of amide group.

tion.^{13,14} More recently, a third helical conformation has been
 described for the benzyl ester of biosynthetic poly(α -benzyl
 γ ,*L*-glutamate).¹² This consists of a 37/10 helix with the same
 hydrogen-bonding scheme and rise per residue (1.50 Å) as
 the 17/5 helix, which make both helices very close in
 topology. In fact, the maximum difference between dihedral
 angles of the 17/5 and 37/10 helices is only 4.1°. Given the
 strong similarities between these two helices and the high
 computing cost that the analysis of the 37/10 helix would
 imply, only the 5/2 and 17/5 helical arrangements were
 considered in this study. However, it should be emphasized
 that results obtained with the 17/5 helix are perfectly
 applicable to the 37/10 helix because of their structural
 similarity. The main characteristics of the 5/2, 17/5, and 37/
 10 helices are summarized in Table 1. Figure 2 shows the
 axial and equatorial projections of 5/2 and 17/5 helices
 examined in the present work.

Molecular Models. The structure of comblike PAADG-
 18 in the solid state was simulated by packing four
 independent helices of PAADG-18 in an orthogonal simula-
 tion box with the helices oriented with their axes parallel to
 the *z*-axis (Figure 3). The *x*-axis was defined as pointing
 along the line joining centers of neighboring helices within
 a layer. The *y*-axis pointed therefore perpendicular to the
 helix layers, and these were parallel to the *x-z* plane.
 Independent chains in the arrangements constituted by 5/2
 and 17/5 helices contained a total of 10 (5 × 2) and 17
 residues in each one, respectively. Periodic boundary condi-
 tions were applied along the three axes by using the
 minimum image convention.

All of the methylene and end methyl groups were
 represented using pseudoatoms, while the remaining atoms
 were described explicitly. The number of explicit atoms/
 pseudoatoms comprising a typical simulation of PAADG-
 18 varied with the helical backbone conformation, being 1160
 and 1972 for the 5/2 and 17/5 helices, respectively. Bond
 lengths and angles were kept fixed for the alkyl side chains,
 whereas bond lengths and angles and torsional angles were
 kept fixed for all helices, which is consistent with the

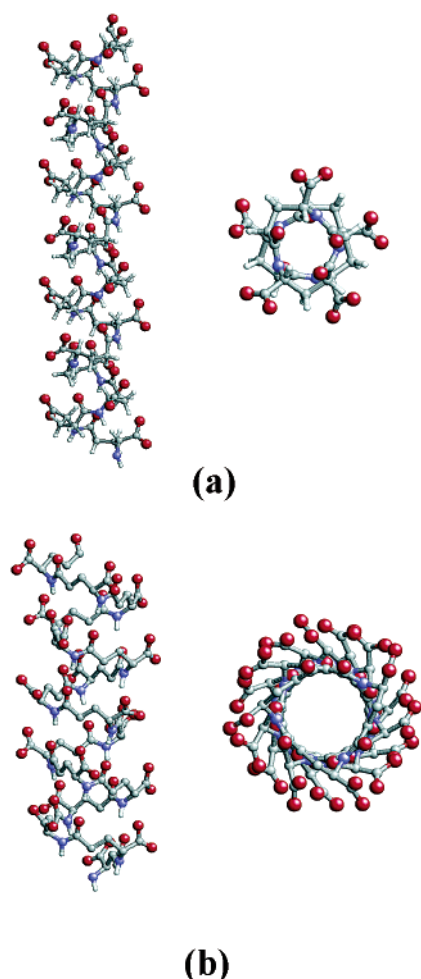


Figure 2. Axial and equatorial projections of (a) the 5/2 and (b) the 17/5 helices considered for PAADG-18. To clarify the figure, the octadecyl side groups have not been represented.

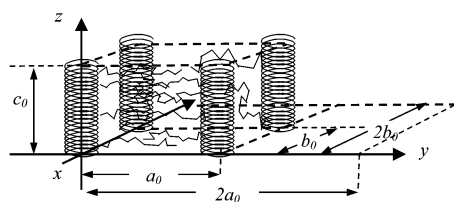


Figure 3. Schematic representation of the simulated models. The cell parameters b_0 and a_0 correspond to the separation between helices within a layer and between two successive layers, respectively. The dimensions of the simulation box for the different models investigated were $2a_0$, $2b_0$, and c_0 .

experimental observation that helical conformations are retained in phases A and B.⁹

Simulation Details. Simulations were performed using an advanced MC sampling technique, configurational bias Monte Carlo (CBMC).^{19,20} The CBMC method consists of the following three steps: (i) a chain is selected at random; (ii) the chain is cut at a random position; (iii) the chain is sequentially regrown bond-by-bond by examining a number of possible torsions (N_s), which are randomly chosen. This algorithm was specifically adapted for studying comblike polymers and subsequently implemented into a computed program denoted MCDP (Monte Carlo Simulations of Dense Polymers).²¹ In addition to CBMC moves, a small fraction of Metropolis moves was also considered for the alkyl side

Table 2. Temperature, Number of Steps, and Frequency (%) for the Different Types of Monte Carlo Moves in *NVT* and *NPT* Simulations

simul	model	T^a	type	steps	CB ^b	Metrop ^b	a/b^c	θ^d
1	5/2-A	348	<i>NVT</i>	7.5×10^4	0.6	0.1		0.3
2	5/2-P	348	<i>NVT</i>	7.5×10^4	0.6	0.1		0.3
3	5/2-A	348	<i>NPT</i>	2.5×10^5	0.55	0.05	0.2	0.2
4	5/2-P	348	<i>NPT</i>	2.5×10^5	0.55	0.05	0.2	0.2
5	17/5-A	348	<i>NVT</i>	7.5×10^4	0.6	0.1		0.3
6	17/5-P	348	<i>NVT</i>	7.5×10^4	0.6	0.1		0.3
7	17/5-A	348	<i>NPT</i>	2.5×10^5	0.55	0.05	0.2	0.2
8	17/5-P	348	<i>NPT</i>	9×10^5	0.55	0.05	0.2	0.2
9	17/5-P	348	<i>NVT</i>	3.5×10^5	0.6	0.1		0.3
10	17/5-P	363	<i>NVT</i>	5×10^5	0.55	0.05		0.4

^a Temperature in K. ^b Frequency of CB and Metropolis moves for the alkyl side chains. ^c Frequency of *NPT* moves. ^d Frequency of moves for the setting angles of the helices.

chains. In the CB algorithm, $N_s = 8$ torsional angles were used to sample the torsional space for the side chains. The degrees of freedom in simulations of *NVT*-type, that is, without varying the size of the simulation box, were the torsional angles of the alkyl side chains and the setting angles, which define the relative orientation among the helices. On the other hand, in *NPT*-type simulations ($P = 1$ atm), the dimensions of the simulation box were also considered as degrees of freedom. The frequency used for the different types of MC moves, the number of steps, and temperature of the system for all of the simulations presented below are displayed in Table 2.

The Amber force field was used to represent the electrostatic, van der Waals, and torsional energies of the system.²² The van der Waals energy was computed in the usual pairwise additive mode using a Lennard-Jones 6-12 potential. The van der Waals parameters, σ and ϵ , were computed using arithmetic and geometric mean combining rules, respectively.

Electrostatic interactions play a major role in PAADG- n helix conformation; they are actually responsible for their high stability. Because we are comparing the stability of the 5/2 and 17/5 helices in supramolecular biphasic structures, these interactions were taken into account by assigning partial atomic charges to the atoms of both the backbone and side ester groups. Such electrostatic charges were derived by fitting the rigorously defined quantum mechanical molecular electrostatic potential, which was calculated at the ab initio HF/6-31G(d) level on a reduced model constituted by two residues, to the Coulombic electrostatic potential. This procedure was used for two reasons. First, previous studies demonstrated that in general the parameters derived from suitable reduced models lead to reliable results.²³ Second, the electrostatic parameters derived at the HF/6-31G(d) level have been included in the libraries of the Amber force fields.²² However, electrostatic interactions were neglected for the mobile alkyl side chains because they can be considered as electrically neutral. This strategy, which is efficient from a computational point of view, was proved to be reliable for comblike PAALA- n .⁸ Thus, calculations including the electrostatic interactions for the alkyl side groups led essentially to the same results as those in which such interactions were omitted. Furthermore, it should be noted that electrostatic interactions are usually neglected in

221 both molecular dynamics and MC simulations of melts of
222 alkyl chains.²⁴

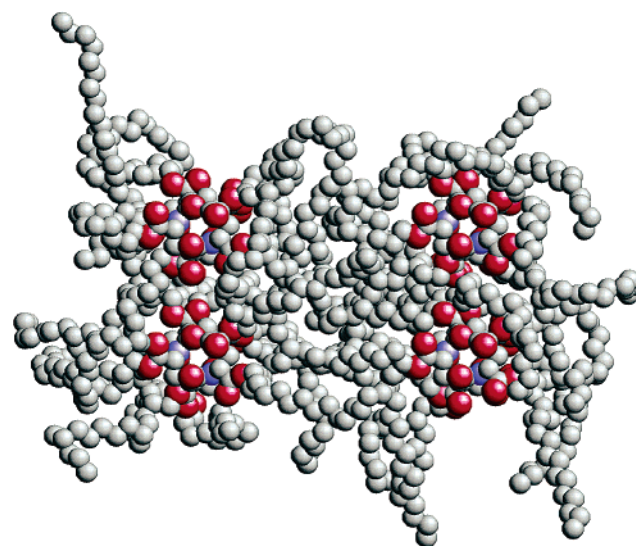
223 Electrostatic interactions for the partially charged atoms
224 were evaluated using a standard Coulombic potential. It is
225 expected that such a simple approach produces reliable results
226 for the following reasons. First, the atomic charges on the
227 backbone atoms are relatively low because no anion/cation
228 is included in the system. Furthermore, the distance between
229 neighboring polymer helices is considerably large. Non-
230 bonding interactions were truncated at 15 Å, implying that
231 all atoms of one residue interact with all other atoms of
232 another residue if at least one pair of atoms is within this
233 limit. A three-term Fourier series expansion was used to
234 represent the torsional energy. Nonbonding and torsional
235 parameters were taken from Amber 4.0 libraries.²⁵

236 It should be mentioned that during the last years several
237 force fields have been optimized for the simulation of long
238 paraffinic chains.²⁶ Within this context, the anisotropic united
239 atom force field developed by Toxvaerd^{26a,b} and co-workers
240 deserves special attention. This model, which allows the
241 movement of the interaction center on each segment depend-
242 ing on the conformation of the whole molecule, was
243 successful in predicting the equation of state and dynamics
244 of alkanes, as well as the structure and thermodynamics of
245 Langmuir monolayers. However, it should be noted that MC
246 simulations of PAADG-18 should provide a satisfactory
247 description not only of the paraffinic interphase but also of
248 the polypeptide chains. The Amber force field is able to
249 satisfy such requirements,⁵⁻⁸ and for this reason, we decided
250 to use it. Furthermore, previous studies indicated that the
251 behavior of long alkyl chains can be also correctly simulated
252 using simple isotropic force fields.^{19,24}

253 The atomistic modeling of the phase A was performed by
254 using the graphical tools implemented in the Cerius 2
255 computer package.²⁷

256 Results and Discussion

257 **The Helical Conformation in Comblike PAADG-18.** For
258 each helix type, two different packing modes were consid-
259 ered: (i) the chains arranged antiparallel with respect to each
260 other (A) and (ii) the chains arranged in parallel (P). The
261 initial dimensions of the simulation box for the two packing
262 modes of 5/2 helices, henceforth denoted 5/2-A and 5/2-P
263 models, were $2a_0 = 56.00$, $2b_0 = 24.21$, and $c_0 = 20.20$ Å
264 (Figure 3). Unfavorable interactions were removed by 7.5
265 $\times 10^4$ MC steps of *NVT*-type at $T = 348$ K (simulations 1
266 and 2 in Table 2), this temperature being about 20 K higher
267 than that observed for the phase A-B thermal transition in
268 PAADG-18.⁹ After this, production runs consisting of 2.5
269 $\times 10^5$ steps of *NPT*-type at the same temperature were started
270 (simulations 3 and 4 in Table 2), the atomic coordinates being
271 saved at 2500 steps intervals. To consider the influence of
272 the starting point in the simulations, additional calculations
273 were performed for the 5/2-A and 5/2-P models by varying
274 both the dimensions of the simulation box and the arrange-
275 ment of the helices (data not shown). However, such
276 simulations do not deserve any extra discussion because they
277 provided similar results to those presented below.



278 **Figure 4.** Equatorial projection (x - y plane) of a representative
279 microstructure provided by MC simulations of *NPT*-type for the 5/2-P
280 model of PAADG-18.

281 Both the 5/2-A and 5/2-P models were found to evolve
282 toward a structure in which the alkyl side chains are
283 completely molten. Inspection of the recorded microstruc-
284 tures indicated that the side chains of adjacent residues were
285 very separated, so they cannot pack favorably. As a
286 consequence, the paraffinic side chains wrap around the
287 helical backbones and large voids appeared in the middle
288 part of the interlayer region. These features are illustrated
289 in Figure 4, which shows a representative microstructure
290 projected along the c -axis for the 5/2-P model.

291 The deficiencies detected for models 5/2-A and 5/2-P must
292 be attributed to the structural characteristics of the 5/2 helix
293 (Table 1). Consecutive side chains are spaced too far apart
294 because of both the large rise per residue (2.02 Å) and the
295 small number of residues per turn (2.5 residues turn⁻¹), which
296 induce their folding toward the polypeptide backbone. The
297 absence of interdigitation among the alkyl side chain of
298 different helices is in contradiction with experimental
299 evidences for PAADG- n ,⁹ and also PGALG- n and PAALA-
300 n .¹⁻⁶

301 Figure 5a shows the evolution of the energy through the
302 *NPT* simulations for the two packing modes constituted by
303 5/2 helices. As can be seen, no significant energy difference
304 appears between the 5/2-A and 5/2-P models, even though
305 the 5/2-P model is slightly more stable (~ 1 kcal/mol) than
306 the 5/2-A. On the other hand, the evolution of the interlayer
307 distance, a_0 , through the simulations is displayed in Figure
308 5b. As can be seen, the interlayer distance a_0 decreases from
309 28 to 25.5 and 24.5 Å for the 5/2-A and 5/2-P, respectively.
310 These values are about 5 Å shorter than the parameter
311 experimentally measured for the phase B of PAADG-18 (a_0
312 = 30 Å).⁹ Regarding the distance between neighboring
313 helices within the same layer, b_0 , it fluctuates around the
314 initial value to arrive to ~ 11.8 Å at the end of the simulation
315 (data not shown). We are aware that the parameter a_0 is not
316 completely equilibrated after 2.5×10^5 MC steps for the
317 5/2-P model. Nevertheless, in our opinion neither the 5/2-A
318 nor the 5/2-P models deserve further consideration. Thus,

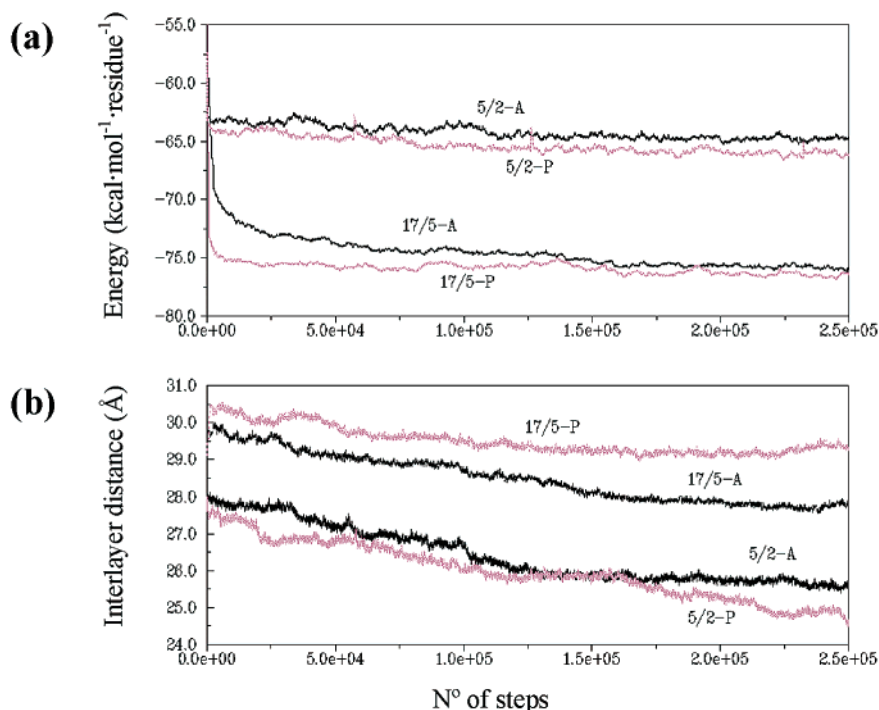


Figure 5. Evolution of (a) the energy and (b) the interlayer distance, a_0 , through the *NPT* MC simulations ($T = 348$ K) for the 5/2-A, 5/2-P, 17/5-A, and 17/5-P models.

316 the overall results allow us to conclude that the models
 317 constituted by 5/2 helices provide a poor description of both
 318 the cell dimensions and the distribution of the alkyl side
 319 chains in the paraffinic interphase.

320 To investigate the second helical conformation considered
 321 in this study, the two packing modes of helices 17/5, denoted
 322 17/5-A and 17/5-P models, were immersed in a simulation
 323 box with dimensions $2a_0 = 58.00$, $2b_0 = 28.60$, and $c_0 =$
 324 25.50 Å. In this case, the parameters $2a_0$ and $2b_0$ were
 325 increased with respect to those initially considered for the
 326 5/2-A and 5/2-P model. Otherwise unfavorable steric clashes
 327 were obtained because the diameter of the 17/5 helix is 1.8
 328 Å larger than that of the 5/2 helix. In a first stage, MC
 329 simulations of *NVT*-type at $T = 348$ K were performed
 330 to minimize unfavorable sterical clashes (simulations 5 and 6
 331 in Table 2). Next, simulations of *NPT*-type were performed
 332 at the same temperature (simulations 7 and 8 in Table 2).

333 It is worth noting that the interlayer distance a_0 sharply
 334 increases during the first hundredths of MC steps (Figure
 335 5b). This enlargement is due to some unfavorable interactions
 336 that remained after *NVT* simulations. Subsequently, a_0
 337 shortens slowly until reaching an equilibrium value. This is
 338 about 28 and 30 Å for the 17/5-A and 17/5-P models,
 339 respectively, revealing a good agreement with the experi-
 340 mental measure ($a_0 = 30$ Å). On the other hand, the
 341 parameter b_0 stabilizes at about 14 Å (data not shown), this
 342 value being similar to that used as starting point. It should
 343 be noted that no experimental value has been reported for
 344 b_0 . The density calculated for this structure is about 1.0
 345 g mL⁻¹, which is in comfortably good agreement with the value
 346 that should be experimentally expected. Note that the density
 347 measured for phase A is 1.02 g mL⁻¹. According to the
 348 contraction happening in the structure when phase A converts
 349 into phase B, the density of the latter would be increased by

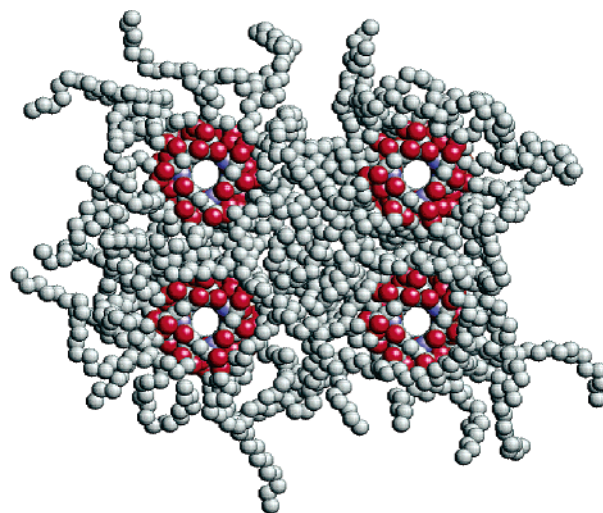


Figure 6. Equatorial projection (x - y plane) of a representative microstructure provided by MC simulations of *NPT*-type for the 17/5-P model of PAADG-18.

6%. Indeed, the small difference (less than 10%) between 350
 the theoretical and experimental densities should be attributed 351
 to the Amber force field, which was optimized to study 352
 biological macromolecules in dilute solution.²² Similar trends 353
 were detected in our previous studies on comblike PAALA- 354
 n.⁶⁻⁸ 355

Figure 6 shows a representative microstructure of the 17/ 356
 5-P model obtained in the production stage. As can be seen, 357
 the appearance of void spaces at the center of the paraffinic 358
 region is minimized because side chains are interdigitated, 359
 which is in good agreement with experimental data.⁹ 360
 Comparison with the results obtained for the 5/2-A and 5/2-P 361
 models indicates that the shape and size of the core defined 362
 by the helix backbone plays a crucial role. Thus, the 17/5 363

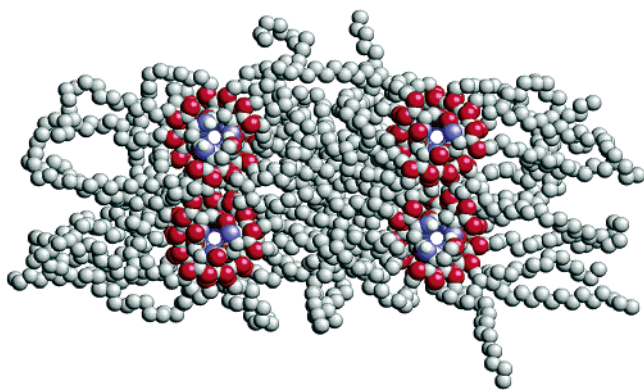


Figure 7. Equatorial projection (x - y plane) of a representative microstructure provided by MC simulations of NPT -type for the phase B of PAALA-18.

helix presents a rise per residue of 1.50 \AA , and its projection can be described as cylindrical. These characteristics allow both favorable interactions among the interdigitated side chains, because they are close in the space, and a homogeneous distribution of the side chains in the bending region, that is, the region in the proximity of the helix backbone.

The evolution of the energy after equilibration of the structures made of $17/5$ helices is shown in Figure 5a. Again, no significant energy difference appears between the parallel and antiparallel models, the $17/5$ -P model being only slightly more favored than the $17/5$ -A one. However, an important energy gap (larger than $12 \text{ kcal mol}^{-1} \text{ residue}^{-1}$) is observed when the models constituted by $17/5$ and $5/2$ helices are compared. The overall results lead us to consider that the $17/5$ helix is more suitable for the formation of the biphasic structures of PAAGD-18 than the $5/2$ one. Furthermore, we considered the $17/5$ -P packing for subsequent analyses even though no clear difference can be established between the parallel and the antiparallel models. For this purpose, simulation 8 was extended to 9×10^5 MC steps, no significant change being found in the lattice dimensions and the energy with respect to the results displayed in Figure 5.

A detailed inspection of the microstructures generated for the $17/5$ models allows detection of a notable disorder in the interlayer region (Figure 6). Accordingly, the phase B of PAADG-18 can be conceived as layers of polypeptide helices embedded in a matrix made up of paraffinic chains in a molten state. This is a striking difference with respect to the phase B of comblike PAALA- n , in which the alkyl side chains retain a preferential alignment along the y -axis, that is, the side chains are partially disordered but not in the actual molten state.^{6,7}

Furthermore, another important difference concerning the side-by-side arrangement of the helices appears between comblike poly(β -peptide)s and poly(γ -peptide)s. In the former case, neighboring helices within a layer were almost in contact. Thus, the alkyl side chains were essentially located in the interlayer region. Figure 7 shows a microstructure of PAALA-18, which was also obtained using MC simulations.⁸ Conversely, the layers of PAADG-18 are constituted by polypeptide helices separated by alkyl side chains in a molten state, that is, the paraffinic chains are distributed between both the inter- and intralayer regions. However, the intralayer paraffinic region is so thin that neighboring helices are able

to interact through their dipoles. Indeed, the cell dimensions derived from MC simulations clearly reflect the existence of such interaction. Thus, the parameters reached after 9×10^5 MC steps follow the pattern displayed in Figure 5b: a_0 ($\sim 30 \text{ \AA}$) $\approx 2b_0$ ($\sim 2 \times 14 \text{ \AA}$). It should be emphasized that in the absence of such interactions the system should spontaneously evolve toward a new isotropic structure with parameters $a_0 \approx b_0$. The latter arrangement is that supposed to exist in the phase C of PAALA-18, which is reached upon heating the phase B upon a second transition temperature T_2 .⁴

The differences between the phase B of PAADG-18 and PAALA-18 should be attributed to the topological differences that derive from the constitutional and conformational characteristics of these compounds. Comblike poly(β -peptide)s adopt a $13/4$ helix with 3.25 residues per turn, while a $17/5$ helix with 3.4 residues per turn is here assumed for PAADG-18. Furthermore, consecutive amide groups are separated along the main chain by two and three carbon atoms in PAALA- n and PAADG- n , respectively. Accordingly, the backbone and side ester atoms of each helix are closer in the former compounds than in the latter ones. This feature explains the presence of alkyl side chains between adjacent helices within the layers and the higher disorder found in the interlayer region.

Comparison with Phase A. The structural conclusions drawn for the phase B concerning the suitability of the $17/5$ helical arrangement should be expected to be even more conspicuous for phase A because in this case the topological restrictions are more severe. Previous studies in comblike PAALA- n indicated that unconstrained MC simulations are not suitable to reproduce the phase A.⁷ The failure to reproduce the crystallization of the paraffinic chains is due to the MC method itself. It is well-known that MC simulations of alkanes below the experimental melting point lead to supercooled liquids rather than to crystals.²⁸

To provide an atomistic model of phase A comparable with phase B modeled by MC methods, we employed the graphical modeling tools implemented in the Cerius2 computer program.²⁷ Thus, one of the microstructures derived in the previous section for the $17/5$ -P model was used as starting point, both the cell dimensions and the torsional angles of the alkyl side chains being adjusted to fulfill the features derived from the experimental data. This process was combined with single-point energy calculations to avoid unfavorable interactions in the resulting model.

Figure 8 shows the atomistic model proposed for phase A of PAADG-18, which is characterized by the presence of order in the tree-axis. Even though this highly ordered structure fits the most important experimental trends reported for the phase A (density and crystallization of about eight methylene groups in the interlayer region), it should be only considered a rough atomistic model. This is because the conformational space of the dihedral angles involved in the side chain bending regions has not been explored. Unfortunately, the limitations of the MC techniques and the lack of experimental data concerning the side group arrangement in such regions does not allow us to obtain a more precise model.

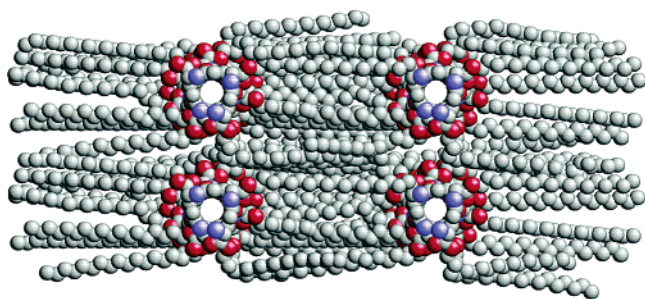


Figure 8. Representative structure of the phase A of PAADG-18 with main chains in the 17/5 helical conformation.

467 **Organization of the Paraffinic Chains in Phase B of**
 468 **PAADG-18.** To provide a more quantitative description of
 469 the structure of the paraffinic region in the phase B of
 470 PAADG-18, the 18 torsional angles for the alkyl side chains
 471 of the 17×4 residues explicitly considered were examined.
 472 For this purpose, additional simulations of *NVT*-type consist-
 473 ing of 3.5×10^5 steps were performed at $T = 348$ and 363
 474 K for the 17/5-P model (simulations 9 and 10 in Table 2).
 475 According to the results presented in previous sections, the
 476 parameters of the simulation box were $2a_0 = 60.0$, $2b_0 =$
 477 28.0 , and $c_0 = 25.5 \text{ \AA}$.

478 The short-distance properties (at a level of $1\text{--}10 \text{ \AA}$) of
 479 the alkyl side chains were examined by considering the first
 480 and second degree autocorrelation functions ($f_{1;\text{bcf}}$ and $f_{2;\text{bcf}}$,
 481 respectively), which allow us to study how fast the local
 482 properties of the paraffinic phase change as the MC simu-
 483 lations run. These were computed using the following expres-
 484 sions:²⁹

$$f_{1;\text{bcf}}(n) = \langle v_i(j)v_i(j+n) \rangle_{ij}$$

$$f_{2;\text{bcf}}(n) = \frac{3}{2} \langle (v_i(j)v_i(j+n))^2 \rangle_{ij} - \frac{1}{2}$$

485 where the subscripts i and j correspond to the bonds and to
 486 the microstructures, respectively, and v_i is the unit vector of
 487 the i th bond. These functions estimate how fast the bonds
 488 erase the memory of the previous local configuration. The
 489 decay of both $f_{1;\text{bcf}}(n)$ and $f_{2;\text{bcf}}(n)$ at $T = 348$ K is shown in
 490 Figure 9, the autocorrelation functions obtained at $T = 363$
 491 K (data not shown) being similar. The bond autocorrelation
 492 functions $f_{1;\text{bcf}}(n)$ and $f_{2;\text{bcf}}(n)$ drop to about 0.7 and 0.6,
 493 respectively. It should be remarked that for the system
 494 investigated, one of the ends of the side chain is anchored
 495 to the helix backbone precluding a complete loss of both
 496 short- and large-range correlations. According to this and to
 497 the results displayed in Figure 9, the microstructures obtained
 498 in the present simulations should be considered as statistically
 499 independent.

500 Figure 10 shows a population analysis for each torsional
 501 angle, the conformations being grouped in the following four
 502 categories: trans, gauche⁺, gauche⁻, and the other remaining
 503 conformers. The trans is the predominant conformation at
 504 the two temperatures, which is consistent with the results
 505 obtained for comblike PAALA- n .⁶⁻⁸ On the other hand, the
 506 frequency of trans conformation slightly decreases with the
 507 temperature, that is, about 5% in average. This is not a
 508 surprising result because the population of folded states is
 509 expected to increase with the temperature.

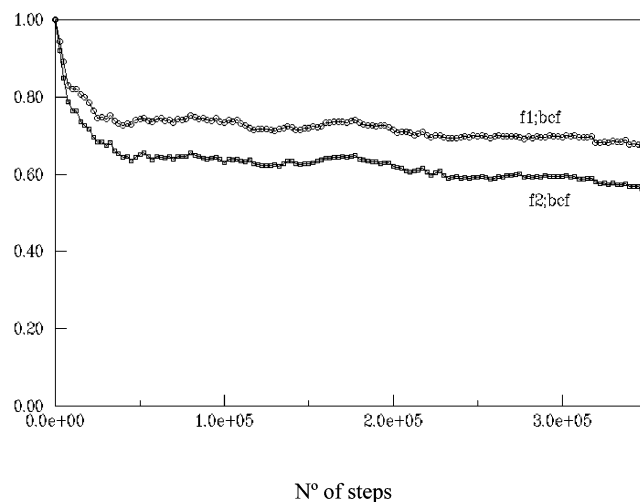


Figure 9. Bond autocorrelation functions ($f_{1;\text{bcf}}$ and $f_{2;\text{bcf}}$) for the phase B of PAAG-18 at $T = 348$ K.

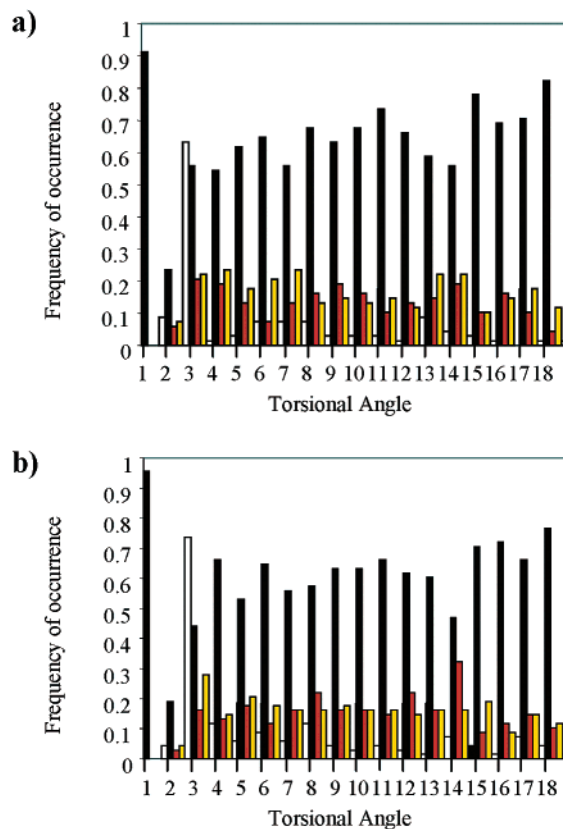


Figure 10. Torsional angle distribution for the alkyl side chains of PAADG-18 at (a) $T = 348$ K and (b) $T = 363$ K (b). The population analysis of the torsional angle associated with each of the 18 bonds in the alkyl side chain is specified. The four categories considered for each bond, in the order displayed in the figure from left to right, are trans, gauche⁺, gauche⁻, and the remaining conformers.

There is a fundamental difference between the results 510
 reported for PAALA-18 and those derived in the present 511
 work for PAADG-18. In comblike PAALA-18, the six 512
 torsional angles closer to the main chain helix present notable 513
 conformational differences with respect to the remaining 12 514
 torsional angles. This feature was related with the alignment 515
 of the alkyl side chains along the y -axis. Conversely, for the 516
 compound under study, no significant difference is perceived 517
 in the populations of the last 16 torsional angles, for which 518

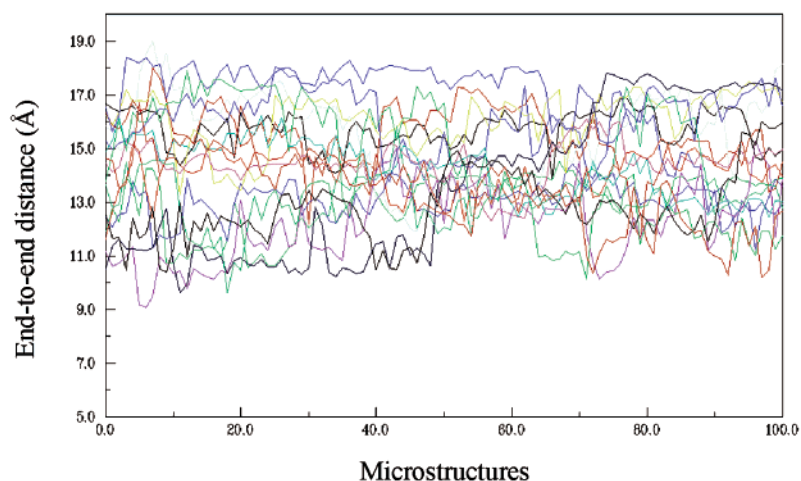


Figure 11. Evolution of the end-to-end distance for the alkyl side chains of PAADG-18 at $T = 348$ K. Every line corresponds to one of the 17 residues explicitly considered.

519 a homogeneous conformational distribution was found. This
 520 is consistent with the disordered state detected in the
 521 paraffinic interphase for the phase B in comblike poly(γ -
 522 peptide)s (see Figure 6).

523 To provide a more detailed view of the homogeneous
 524 distribution of the conformational preferences, the distance
 525 between the carbon atom of the ester group and the end
 526 methyl pseudoatom of the alkyl side chain, that is, the end-
 527 to-end distance of the paraffinic chains, was measured for
 528 each residue. Figure 11 shows the evolution of such distance
 529 for the 17 residues of the helix by considering the last 100
 530 microstructures generated in the simulations at $T = 348$ K.

531 The computed values change with the position of residue
 532 within the layer varying from 11 to 18.5 Å. However, in
 533 general, the predominant values are those ranging from 13
 534 to 16 Å, which indicates that, on average, the different side
 535 chains present similar conformational preferences. Moreover,
 536 comparison between the end-to-end distances at 348 and 363
 537 K (data not shown) reveals that, in general, this parameter
 538 tends to decrease when the temperature increases. This
 539 behavior is in agreement with the population analyses
 540 displayed in Figure 10.

541 Summary

542 The structure of comblike PAADG-18 has been investi-
 543 gated using atomistic Monte Carlo simulations. A standard
 544 nonoptimized force field has been used together with a fixed
 545 bond length, fixed bond angle model. The CB algorithm,
 546 which was specially adapted to simulate comblike polymers,
 547 has been used through the MCDP computer program.²¹

548 Results indicate that the most favored model corresponds
 549 to a parallel packing of 17/5 helices. This helical conforma-
 550 tion is the model previously put forward for poly(γ ,L-
 551 glutamic acid), which is very similar to the 13/4 and α -helix
 552 proposed for PAALA- n and PGALG- n , respectively. The
 553 model would be perfectly replaceable by the 37/10 helix,
 554 the helix that has been found to exist in poly(α -benzyl γ ,DL-
 555 glutamate). On the other hand, the organization of the
 556 paraffinic phase was examined by analyzing both the dihedral
 557 angles of the alkyl side chains and the end-to-end distance.
 558 The structure of comblike PAADG- n can be envisaged

559 therefore as layers of rigid 17/5 helices separated by a molten
 560 paraffinic interphase. The differences detected between the
 561 paraffinic interphase of comblike PAADG- n and that found
 562 for PAALA- n can be considered one of the outstanding
 563 results of this paper.

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